

## TD n°6 : Out of equilibrium statistical physics Kramers-Kronig relation and causality

### 1 Kramers-Kronig relations

We denote by  $\chi(\omega)$  the Fourier transform of a causal function, assumed summable  $\int_{-\infty}^{+\infty} d\omega |\chi(\omega)| < \infty$ . We recall the Kramers-Kronig relations :

$$\operatorname{Re} \chi(\omega) = \frac{1}{\pi} \mathcal{P} \int_{-\infty}^{+\infty} d\omega' \frac{\operatorname{Im} \chi(\omega')}{\omega' - \omega} \quad (1)$$

$$\operatorname{Im} \chi(\omega) = -\frac{1}{\pi} \mathcal{P} \int_{-\infty}^{+\infty} d\omega' \frac{\operatorname{Re} \chi(\omega')}{\omega' - \omega} \quad (2)$$

Where the Cauchy principle value is defined by

$$\mathcal{P} \int_{-A}^{+B} dx \frac{f(x)}{x} \equiv \int_{-A}^{+B} dx f(x) \mathcal{P} \frac{1}{x} = \lim_{\epsilon \rightarrow 0^+} \left( \int_{-A}^{-\epsilon} + \int_{+\epsilon}^{+B} \right) dx \frac{f(x)}{x} \quad (3)$$

where  $A, B > 0$ . Compute the integral for  $f(x) = 1$ .

1/ Given  $\operatorname{Im} \chi(\omega) = -\frac{1}{1+\omega^2}$ . Deduce  $\operatorname{Re} \chi(\omega)$  and  $\chi(\omega)$ .

Hint : In order to compute the Hilbert transform of  $\operatorname{Im} \chi(\omega)$ , integrate  $f(z) = \frac{1}{(z-\omega)(1+z^2)}$  over an appropriate contour in the complex plane.

2/ Same question for  $\operatorname{Im} \chi(\omega) = \frac{\gamma\omega}{(\omega_0^2 - \omega^2)^2 + \gamma^2\omega^2}$ .

### 2 Response function in a deterministic case : harmonic and anharmonic oscillators

Reminder : We denote by  $B(t)$  a physical observable (ex : the position of a particle). We introduce an external time dependent “force”  $f(t)$  coupled to another observable  $A(t)$ , i.e. energy is  $H_{\text{ext}}(t) = -f(t)A$ . The evolution of the observable  $B$  can be linearised as

$$B_f(t) = B_{f=0}(t) + \int_{-\infty}^{+\infty} dt' \chi_{BA}(t-t') f(t') + O(f^2) \quad (4)$$

where  $\chi_{BA}(t)$  is the response function.

1/ **Harmonic oscillator.**— We consider the harmonic oscillator described by the equation of motion  $\ddot{x} + \omega_0^2 x = \frac{1}{m} f(t)$ . Show that the response function  $\chi(t)$  characterising the response of  $x(t)$  to the force  $f(t)$  is the Green’s function of the differential equation. Check that the causal Green’s function is  $\chi(t) = \theta_H(t) \frac{\sin \omega_0 t}{m\omega_0}$ . Compute its Fourier transform  $\tilde{\chi}(\omega)$  (for this purpose it is necessary to introduce a regulator  $e^{-\epsilon t}$  with  $\epsilon \rightarrow 0^+$  in the integral). Plot neatly  $\tilde{\chi}(\omega)$ .

2/ **Damped harmonic oscillator.**— We consider now a damped harmonic oscillator submitted to the external force :

$$\ddot{x} + \frac{2}{\tau} \dot{x} + \omega_0^2 x = \frac{1}{m} f(t) \quad (5)$$

Compute the Fourier transform of the response function  $\tilde{\chi}(\omega)$ . Analyse the poles of this function (for the various regimes). Interpret their positions. Plot neatly  $\text{Re } \tilde{\chi}(\omega)$  and  $\text{Im } \tilde{\chi}(\omega)$  in the weak damping limit (to be defined). Come back to the first question and interpret physically the regulator  $\epsilon \rightarrow 0^+$ .

**3/ Anharmonic oscillator.**— We now consider the classical anharmonic oscillator described by the equation of motion  $\ddot{x} = F(x)$ , where  $F(x)$  derives from a confining potential (e.g.  $V(x) = \frac{1}{2}\omega^2 x^2 + \frac{1}{4}\lambda x^4$ ). Deduce the differential equation satisfied by the response function characterizing the out-of-equilibrium situation  $\ddot{x} - F(x) = f(t)$ . Discuss the differences with the harmonic case.