

## TD n°8 : Out of equilibrium statistical physics (quantum) linear response theory

### 1 Quantum harmonic oscillator

We study the dynamics of a quantum harmonic oscillator described by the Hamiltonian  $H_0 = \frac{p^2}{2m} + \frac{1}{2}m\omega_0^2 x^2$ , submitted to the time dependent perturbation  $H_{\text{pert}}(t) = -x f(t)$ .

1/ Determine the time dependence of the annihilation operator in the interaction representation  $a(t) \stackrel{\text{def}}{=} e^{iH_0 t/\hbar} a e^{-iH_0 t/\hbar}$ .

2/ Deduce the correlation function  $C(t) \equiv C_{xx}(t) = \langle x(t)x \rangle$ , where  $\langle \dots \rangle$  denotes the canonical averaging. Analyze the classical ( $\hbar \rightarrow 0$ ) and zero temperature limits.

3/ Check that the detailed balance relation  $\tilde{C}(-\omega) = \tilde{C}(\omega)e^{-\beta\hbar\omega}$  holds. Deduce the symmetrized correlation function  $S(t) \stackrel{\text{def}}{=} \frac{1}{2}\langle \{x(t), x\} \rangle$ .

4/ Compute the spectral function  $\xi(t) \stackrel{\text{def}}{=} \frac{1}{2\hbar}\langle [x(t), x] \rangle$  and the response function  $\chi(t) \equiv \chi_{xx}(t)$ . Are these results dependent on the nature of the average  $\langle \dots \rangle$ ? Compute  $\tilde{\chi}(\omega)$  and  $\tilde{\xi}(\omega)$ .

5/ Compute the Kubo correlation function  $K(t) \equiv K_{xx}(t) = \frac{1}{\beta} \int_0^\beta d\lambda \langle x(-i\hbar\lambda)x(t) \rangle$ . Check the relation with the response function.

6/ Verify that the response function coincides with the classical result. Give the origin of this observation.

Hint : Write the equation of motion in *Heisenberg* representation for  $x_H(t)$  and  $p_H(t)$ . Justify that the response of the quantum harmonic oscillator is purely *linear*.

7/ Check that the correlation functions satisfy the Fluctuation-Dissipation theorem.

### 2 Conductivity of the electron gas

We consider a gas of (non interacting) free electrons and study its response to an electric field. Perturbation takes the form  $H_{\text{pert}}(t) = -e\vec{\mathcal{E}}(t) \cdot \vec{r}$  (where  $\vec{r}$  is the position operator for *one* electron). The spatially averaged current density is  $\vec{j} = \frac{e}{V}\vec{v}$  where  $\vec{v}$  is the speed operator for the electron. Conductivity (for  $\vec{q}=0$ ) is defined by  $\langle j_i(t) \rangle_{\mathcal{E}} = \int dt' \sigma_{ij}(t-t')\mathcal{E}_j(t')$  (with implicit summation over repeated index).

1/ Compute the conductivity  $\sigma_{ij}^{(1e^-)}(t)$  for a single electron.

2/ Deduce the conductivity for the Fermi gas of  $N$  electrons at temperature  $T$ .

3/ Compute the Fourier transform  $\tilde{\sigma}_{ij}(\omega)$ . Interpret physically the  $\omega \rightarrow 0$  behaviour.

4/ Sum rule.– Give  $\sigma_{ij}(t=0)$  and then  $\int d\omega \tilde{\sigma}_{ij}(\omega)$ .

### 3 Response function for identical particles

We introduce the basis of one particle stationary states  $\{|\varphi_\alpha\rangle\}$ . We denote  $h$  the one body Hamiltonian and  $H = \sum_{i=1}^N h^{(i)}$  the many body Hamiltonian, acting in the  $N$  particle Hilbert space  $\mathcal{H}_N = \mathcal{H}_1^{\otimes N}$ . We denote by  $a_\alpha$  and  $a_\alpha^\dagger$  the annihilation/creation operators acting in the Fock space  $\mathcal{F} = \bigoplus_{N=0}^\infty \mathcal{H}_N$ . They obey  $\{a_\alpha, a_\beta^\dagger\} = \delta_{\alpha,\beta}$  for fermions and  $[a_\alpha, a_\beta^\dagger] = \delta_{\alpha,\beta}$  for

bosons. We consider two observables  $Q$  and  $P$ , sum of one body operators,  $Q = \sum_i q^{(i)}$  and  $P = \sum_i p^{(i)}$ . We recall that, in the Fock space, these operators can be represented as

$$Q = \sum_{\alpha,\beta} q_{\alpha\beta} a_{\alpha}^{\dagger} a_{\beta} \quad (1)$$

where  $q_{\alpha\beta} \stackrel{\text{def}}{=} \langle \varphi_{\alpha} | q | \varphi_{\beta} \rangle$  is the matrix element of the one body operator.

We introduce the grand canonical averaging

$$\langle \dots \rangle \stackrel{\text{def}}{=} \text{Tr} \{ \rho \dots \} \quad \text{with } \rho = \frac{1}{\mathcal{Z}} e^{-\beta(H-\mu N)} \quad (2)$$

where  $\text{Tr} \{ \dots \}$  is a trace in  $\mathcal{F}$  and  $N$  the particle number operator.

1/ Compute  $\langle a_{\alpha}^{\dagger} a_{\beta} \rangle$ . We introduce the notation  $f_{\alpha} \equiv f(\varepsilon_{\alpha}) = \frac{1}{e^{\beta(\varepsilon_{\alpha}-\mu)} \pm 1}$  where  $h | \varphi_{\alpha} \rangle = \varepsilon_{\alpha} | \varphi_{\alpha} \rangle$ .

2/ Verify

$$[AB, CD] = A[B, C]D + [A, C]BD + CA[B, D] + C[A, D]B \quad (3)$$

$$[AB, CD] = A\{B, C\}D - \{A, C\}BD + CA\{B, D\} - C\{A, D\}B \quad (4)$$

and deduce  $[a_{\alpha}^{\dagger} a_{\beta}, a_{\mu}^{\dagger} a_{\nu}]$ .

3/ Show that the following relation

$$\boxed{\langle [P, Q] \rangle = \text{Tr} \{ \rho [P, Q] \} = \text{tr} \{ f(h)[p, q] \}} \quad (5)$$

holds both for bosons *and* fermions, where  $\text{tr} \{ \dots \}$  is a trace in the one body Hilbert space  $\mathcal{H}_1$ .

4/ Note that  $\langle PQ \rangle \neq \text{tr} \{ f(h)pq \}$ .

5/ **IMPORTANT** : Deduce the spectral representation for the response function  $\tilde{\chi}_{PQ}(\omega)$  in terms of one-body matrix elements  $q_{\alpha\beta}$  and  $p_{\alpha\beta}$ . Conclusion ?

## 4 Compressibility of the free fermionic gas

The operator “density of particle” is denoted  $\hat{n}(r)$  and its Fourier component  $\hat{n}_q$  (I drop the arrows on vectors). We consider that the system is in a box of finite size  $V$  and study its response to an external scalar potential  $V^{\text{ext}}(r, t)$ .

1/ Density-density response function  $\chi(r, t; r', 0)$  is defined by

$$\langle \hat{n}(r, t) \rangle_{V^{\text{ext}}} = \overbrace{\hat{n}}^{\text{mean density}} + \int dr' dt' \chi(r, t; r', t') V^{\text{ext}}(r', t') + \dots \quad (6)$$

Express  $\hat{H}_{\text{pert}}(t)$  in terms of  $\hat{n}_q$ . Deduce  $\int d(r - r') e^{-iq(r-r')} \chi(r, t; r', 0)$  as a correlator of the density operator  $\hat{n}_q$ .

2/ We now consider the free fermion gas. One particle eigenstates are the plane waves  $\varphi_k(r) = \frac{1}{\sqrt{V}} e^{ikr}$  with energy  $\varepsilon_k$ . For one fermion, the density operator is  $\hat{n}(r) = \delta(r - \hat{r})$ . Compute the matrix element  $\langle \varphi_k | \hat{n}_q | \varphi_{k'} \rangle$ .

3/ Show that  $\tilde{\chi}(q, \omega) = \int d(r - r') e^{-iq(r-r')} \int dt e^{i\omega t} \chi(r, t; r', 0)$  can be represented as a sum over one-particle eigenstates

$$\tilde{\chi}(q, \omega) = \frac{1}{V} \sum_k \frac{f_k - f_{k+q}}{\hbar\omega + \varepsilon_k - \varepsilon_{k+q} + i0^+} \quad (7)$$

where  $f_k \equiv f(\varepsilon_k)$  is the Fermi-Dirac distribution.

4/ *Static case*  $\omega = 0$ . – Compute  $\tilde{\chi}(q \ll k_F, 0)$ . Interpret the result physically.