

# Anomalous Diffusion of Magnetic Monopoles in Spin Ice

• mpi<mark>pks</mark>

Emergence of a Dynamical Fractal in a Clean Magnet

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#### Spin Ice Basics



Ising spins on the pyrochlore lattice

Constrained to point along easy axes (in or out)

Ferromagnetic nearest-neighbour interactions



→ Frustrated magnet with highly degenerate 2-in-2-out groundstates

**Ice Rules** 

2 spins in and 2 spins out per tetrahedron

Bramwell & Harris, J. Phys.: Cond. Mat. 32 (2020).

# Dysprosium Titanate (Dy<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub>)

J = 15/2 spins with Ising-like single ion states

Long-ranged dipolar spin-spin interactions

Follows the ice rules due to screening



# Dysprosium Titanate (Dy<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub>)

J = 15/2 spins with Ising-like single ion states

Long-ranged dipolar spin-spin interactions

Follows the ice rules due to screening

Emergent magnetic monopole excitations

Monopoles live on the diamond lattice

















Monopole creation cost  $\sim 4$  K Monopole movement cost  $\sim \pm 0.05$  K





### "Standard Model" (SM) of Spin Ice Dynamics

#### Spin flip dynamics (T < 10 K)

Quantum tunnelling between Ising states, enabled by transverse fields.

Flip attempts at constant rate  $1/\tau_0$ .

Monte Carlo time  $\propto$  real time.

Monopole creation is rare at low T.

Monopole motion dominates dynamics.



#### Random choice between 3 directions.

Ryzhkin, J. Theor. and Exp. Phys., 128, 559 (2005). Jaubert & Holdsworth, Nature Physics 5, 258 (2009).

#### The Puzzles



Previous explanations invoked extrinsic contributions (e.g. disorder, boundary effects).

**Anomalous Magnetic Noise** 



Experimental results from: A. M. Samarakoon, et al., Proceedings of the National Academy of Sciences 119, e2117453119 (2022).

#### Beyond the "Standard Model" (bSM)



 $\sim \frac{1}{3}$  of spins neighbouring monopole have  $\vec{B}_{\text{transverse}} = \vec{0}$ .

These attempt flips at lower rate  $1/\tau_{slow}$ .

$$Dy_2 Ti_2 O_7: \qquad \frac{\tau_{slow}}{\tau_{fast}} \approx 10^4$$

(We approximate  $\tau_{slow} = \infty$ )

#### The Puzzles Revisited

#### **Rapidly Diverging Relaxation Time** 10<sup>0</sup> 10<sup>2</sup> Experimental SM SM bSM Relaxation Time (s) 10<sup>0</sup> bSM $10^{-1}$ (in 10<sup>-2</sup> (in 10<sup>-2</sup>) (in 10<sup>-4</sup>) (in 10<sup>-4</sup>) $10^{-2}$ $10^{-6}$ T = 1.04 KT = 0.84 K10-3 $10^{-8}$ T = 0.65 K0.8 1.0 1.2 1.6 $10^{-2}$ 10<sup>0</sup> 10<sup>2</sup> 10<sup>3</sup> 0.6 1.4 1.8 $10^{1}$ 104 10<sup>5</sup> $10^{-1}$ Temperature (K) Frequency (Hz)

Rapidly diverging relaxation time and anomalous magnetic noise explained through purely intrinsic effects!

SM:  $\tau_0 = 200 \ \mu s$ Fitting parameter bSM:  $\tau_{\text{fast}} = 85 \ \mu \text{s}$ 

#### **Anomalous Magnetic Noise**

### Aside: Percolation Theory

Pick your lattice of choice

Randomly fill  $p \ (0 \le p \le 1)$  of the bonds

A single percolating cluster appears at the critical point  $p = p_c$ 

The percolating cluster is self-similar → Fractal



## Links to Percolation Theory

Ice rules and slow spins leave on average 2 directions for a monopole to move in.

- $\rightarrow$ (Dynamical) bond percolation problem on the diamond lattice.
  - $\rightarrow$  Monopoles are random walkers on percolation clusters.

Close to critical filling fraction  $p_c \approx 0.39 \rightarrow$  Fractal structure on length scales up to correlation length.

→ Monopoles move on an **emergent dynamical fractal**!





n = 3

6



n = 3





#### Cluster Growth



#### Monopole Noise



## Subdiffusive Monopoles



#### Summary

Bimodal distribution of internal transverse fields proves crucial to spin ice dynamics.





Explains dynamical properties of spin ice as a consequence of *intrinsic* effects.

Emergent fractal structure in a uniform, disorder-free bulk magnetic crystal.



# Extra slides



#### Hamiltonians

$$\begin{aligned} \mathcal{H}_{\rm NN} &= -J_{\rm eff} \sum_{\langle i,j \rangle} \boldsymbol{S}_i \cdot \boldsymbol{S}_j \\ \mathcal{H} &= D a^3 \sum_{i < j} \left[ \frac{\boldsymbol{S}_i \cdot \boldsymbol{S}_j}{r_{ij}^3} - \frac{3(\boldsymbol{S}_i \cdot \boldsymbol{r}_{ij})(\boldsymbol{S}_j \cdot \boldsymbol{r}_{ij})}{r_{ij}^5} \right] + J_1 \sum_{\langle i,j \rangle} \boldsymbol{S}_i \cdot \boldsymbol{S}_j \\ &+ J_2 \sum_{\langle i,j \rangle_2} \boldsymbol{S}_i \cdot \boldsymbol{S}_j + J_3 \sum_{\langle i,j \rangle_3} \boldsymbol{S}_i \cdot \boldsymbol{S}_j + J_3' \sum_{\langle i,j \rangle_3'} \boldsymbol{S}_i \cdot \boldsymbol{S}_j \\ D a^3 &= 1.3224 \text{ K} \quad J_1 = 3.41 \text{ K} \quad J_{\rm eff} = J_{\rm eff}(T) \sim 5.7 \text{ K} \\ J_2 &= 0.0 \text{ K} \quad J_3 = -0.00466 \text{ K} \quad J_3' = 0.0439 \text{ K} \end{aligned}$$

 $J_2$  $J_3$  $J'_3$ 

#### Cluster Growth



#### Structured percolation



#### Monopole Noise



#### Structured percolation



#### Loop updates



#### Varying slow timescale



#### **Relaxation Time**

