

# High- $T_c$ superconductivity in 2D ruthenates: Relation to CDW/SDW

Sci. Rep. 10, 3462 (2020)

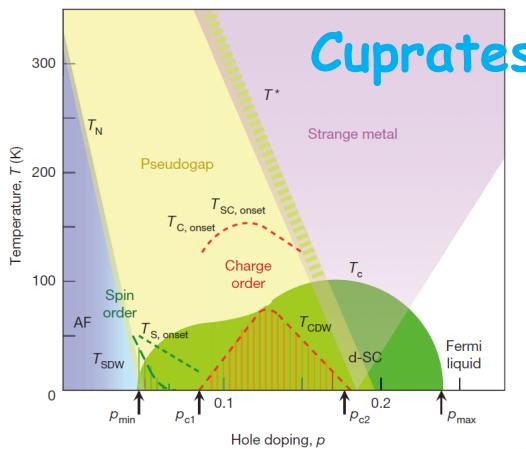
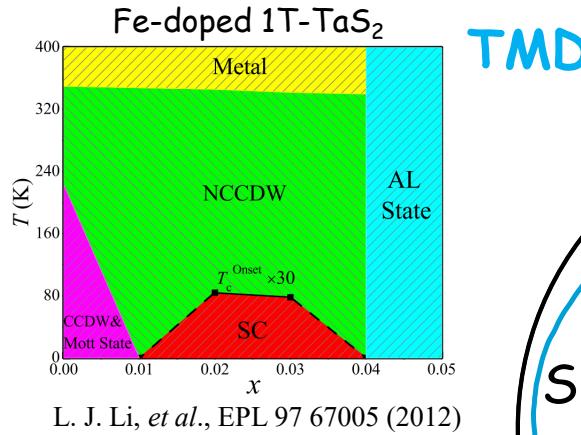
Hiroyoshi Nobukane

K. Yanagihara, Y. Kunisada, Y. Ogasawara, K. Isono, K. Nomura,  
K. Tanahashi, T. Nomura and S. Tanda

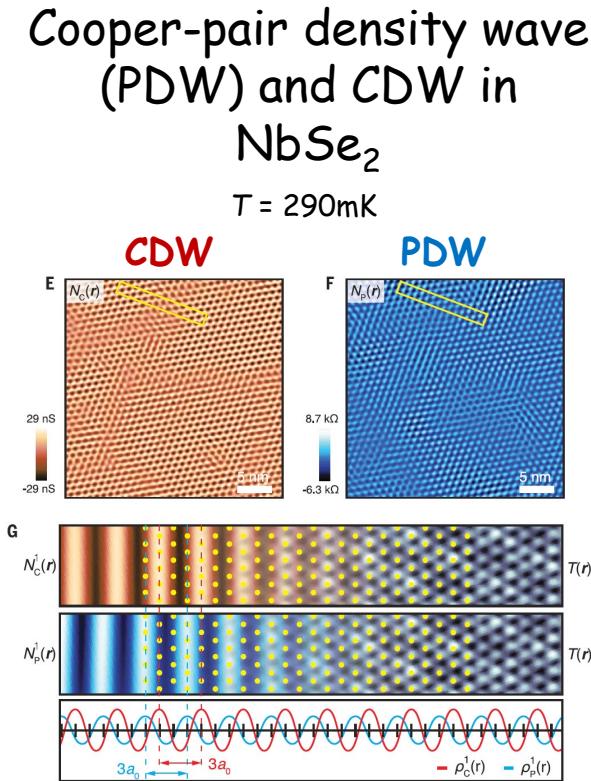
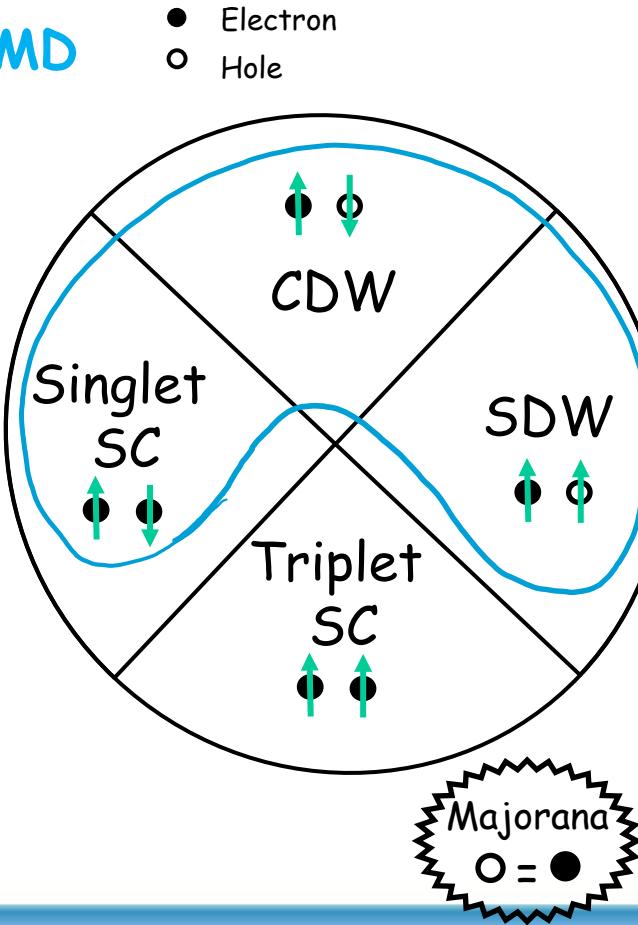


[nobukane@sci.hokudai.ac.jp](mailto:nobukane@sci.hokudai.ac.jp)

# 2D layered materials: TMD and cuprates

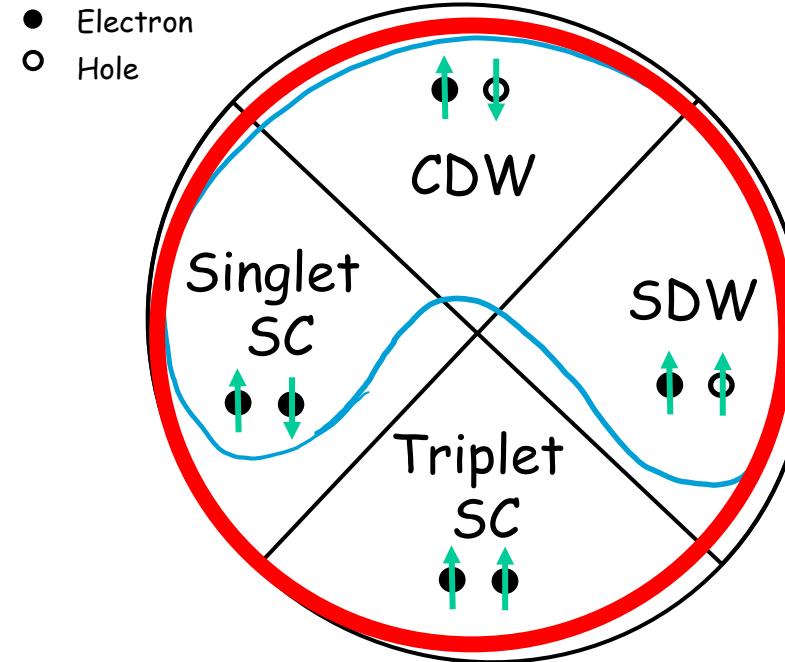


B. Keimer et al., Nature (2015)



X. Liu et al., Science 372, 1447 (2021)

# 2D materials: Triplet SC and DWs !?

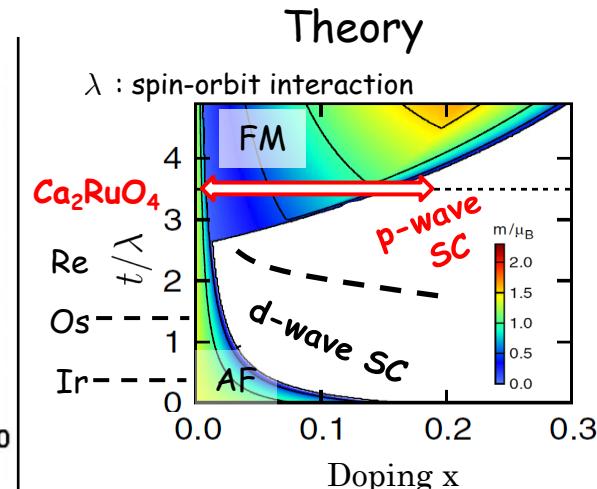
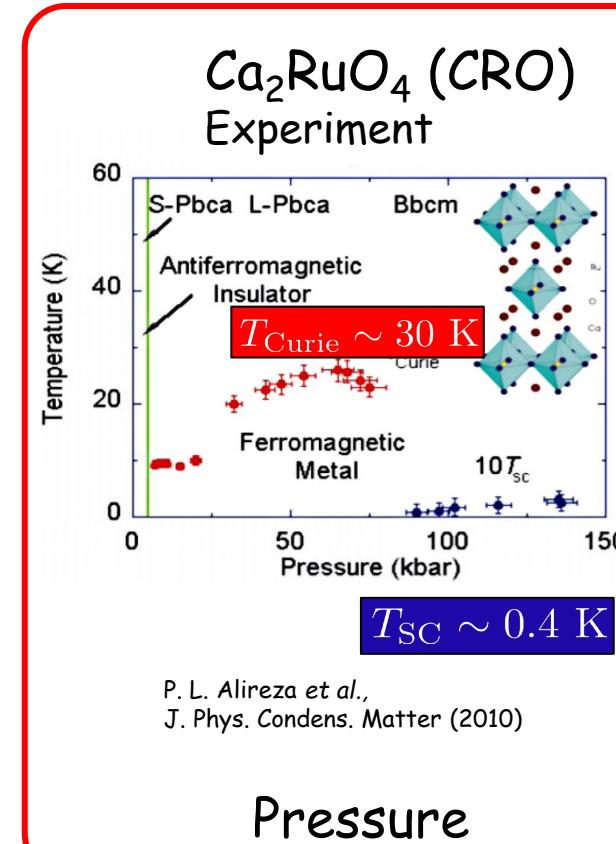
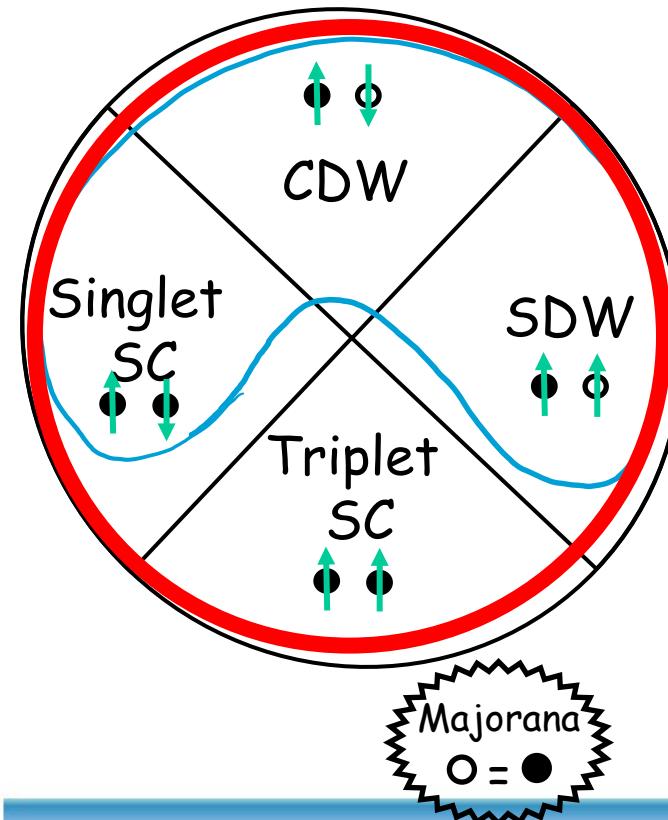


Are there any two-dimensional materials  
in which triplet SC and DWs can be studied?

# 2D ruthenates: SC and FM

- Electron
- Hole

## Ruthenates



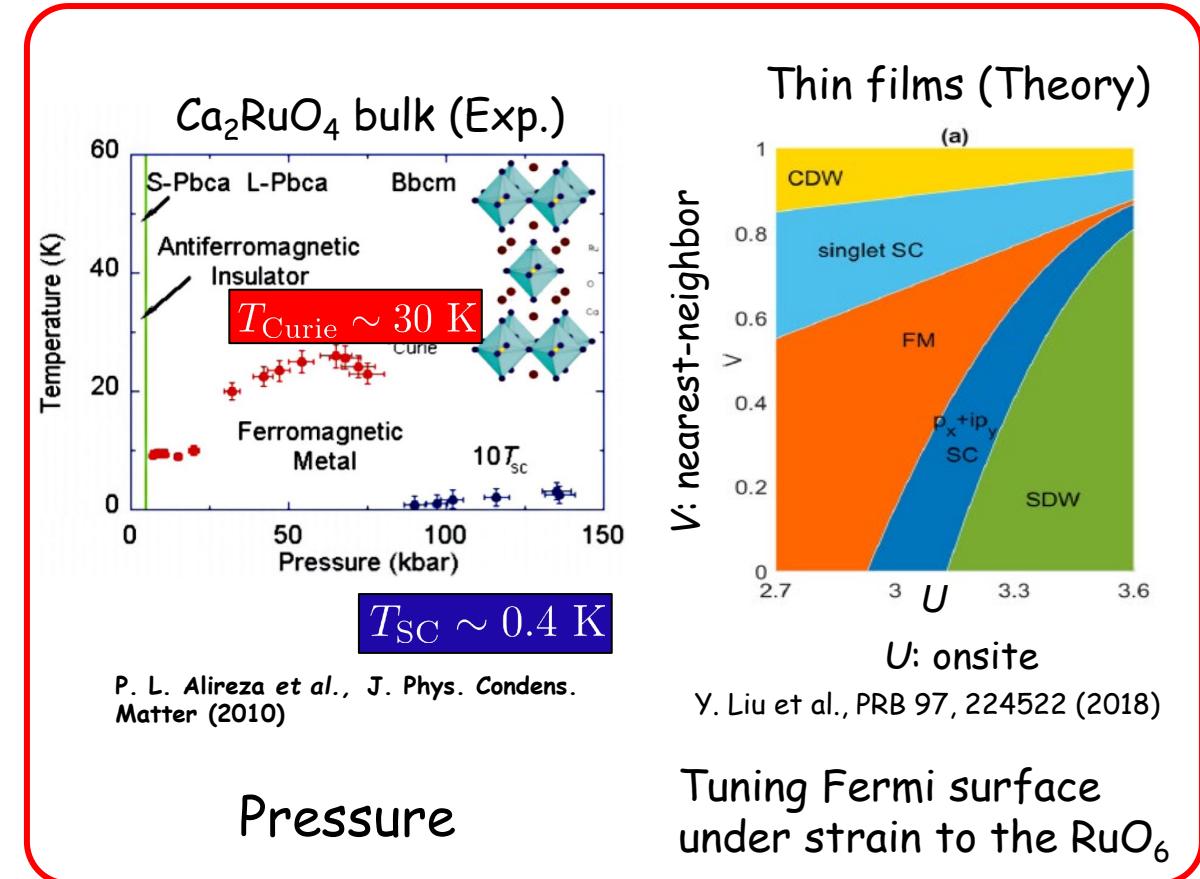
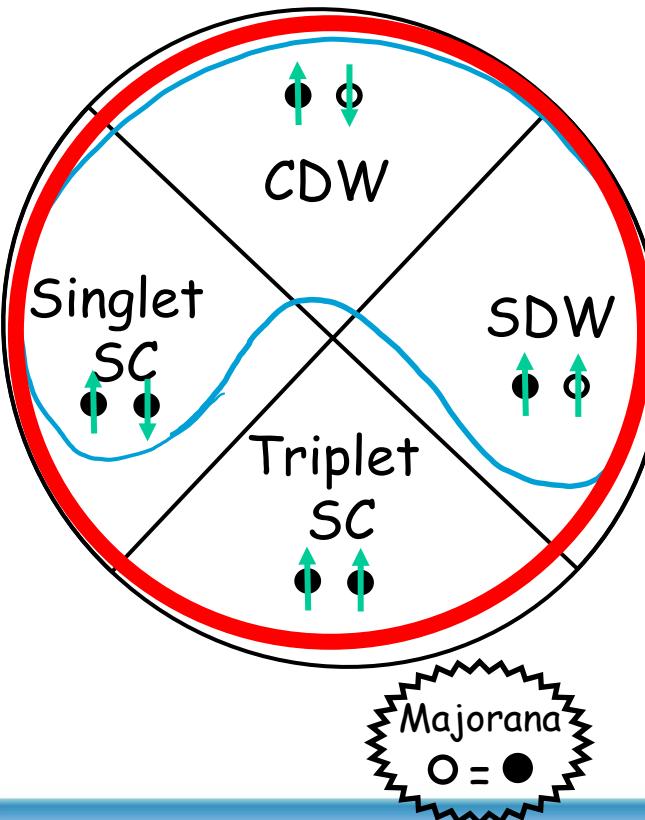
Pressure

Carrier doping

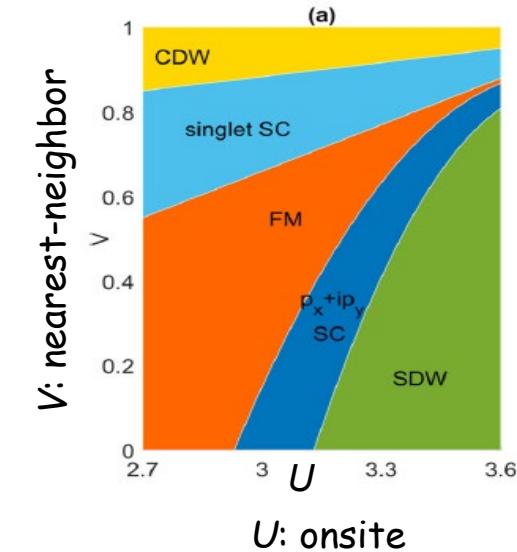
# 2D ruthenates: DWs

- Electron
- Hole

Ruthenates



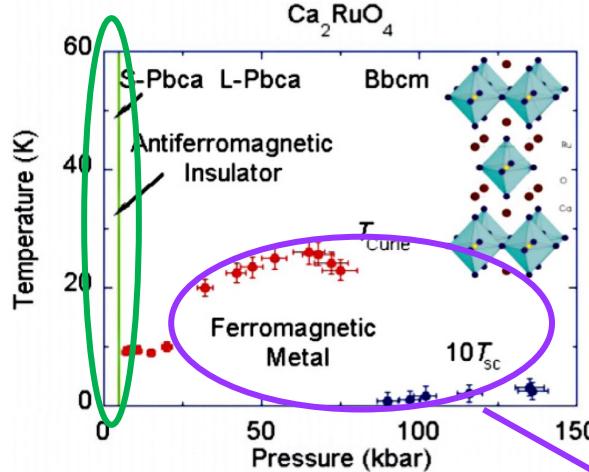
Thin films (Theory)



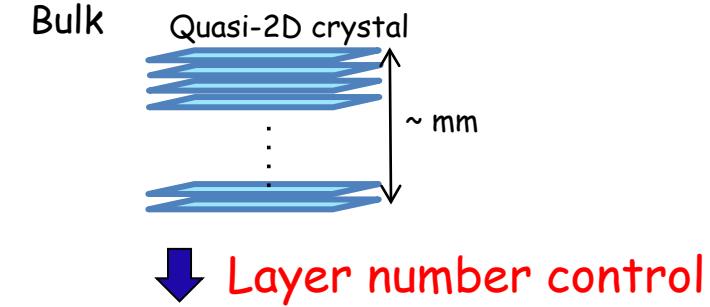
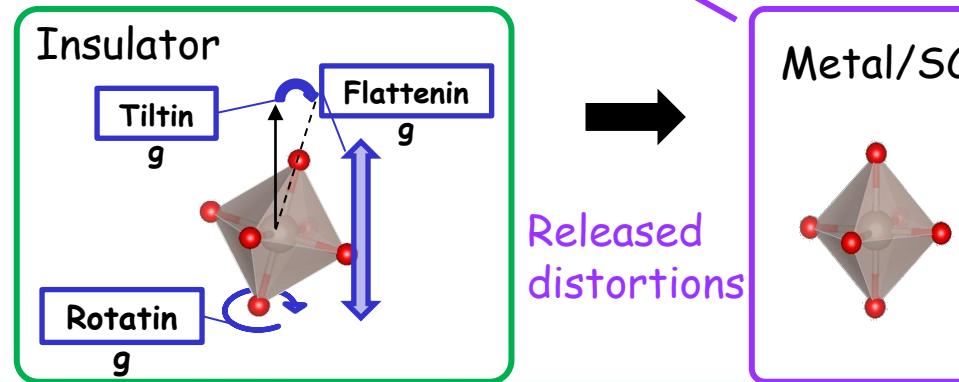
Y. Liu *et al.*, PRB 97, 224522 (2018)

Tuning Fermi surface under strain to the RuO<sub>6</sub>

# Layer number control in ruthenates



P. L. Alireza et al., J. Phys. Condens. Matter (2010)

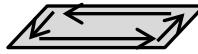


2D ruthenates

**The emergence of novel quantum states!**  
Triplet pairing,  
CDW/SDW

# $\text{Sr}_2\text{RuO}_4$ thin films

1D Edge



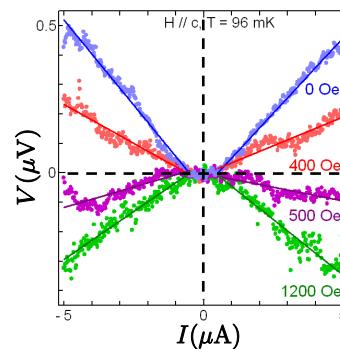
2D Hall



3D Bulk

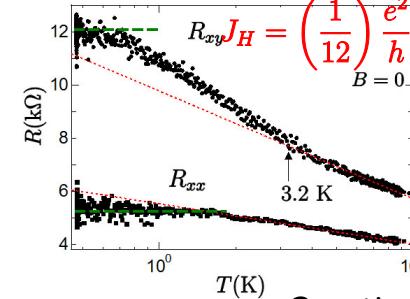
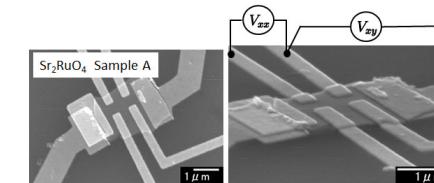


Parity violation of  $I-V$

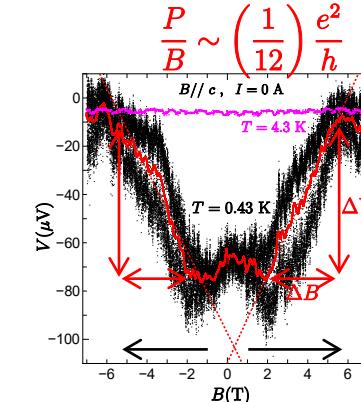


**Chiral Majorana fermions**

- H. Nobukane et al., PRB **83**, 144502 (2011)
- H. Nobukane et al., JJAP **49**, 020209 (2010)
- H. Nobukane et al., SSC **149**, 1212 (2009)



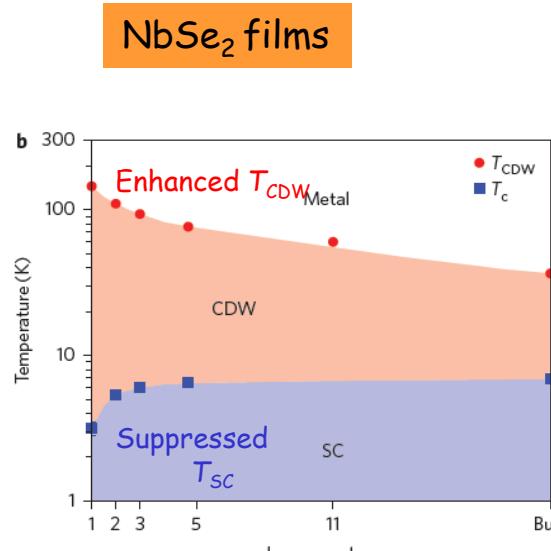
Quantized Hall



**Topological Magneto-electric effect**

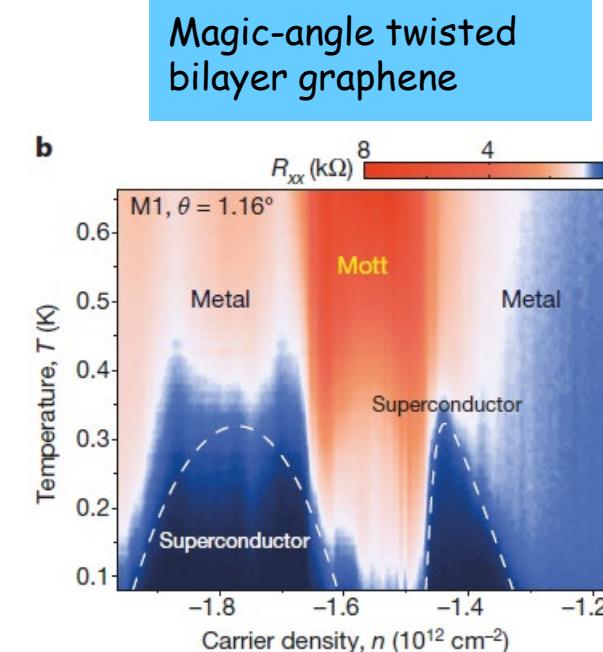
- H. Nobukane et al., Sci. Rep. **7**, 41291 (2017)

# Layer number control in $\text{NbSe}_2$ and graphene



X. Xi *et al.*, *n.nano* **10**, 765 (2015)

Exotic states  
due to a negative pressure!



Y. Cao *et al.*, *Nature* **556**, 43 (2018)

2D layer play a key part for studying ruthenate physics.

## Our purpose

To study electric states from Mott insulator to superconductivity in  $\text{Ca}_2\text{RuO}_4$  nanocrystals by reducing the number of layers

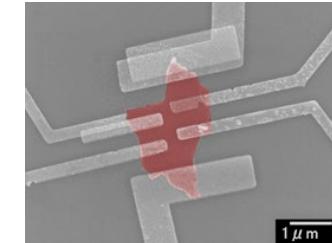
## Experimental

### Nano-crystal

Synthesized nanoscale crystals  
with a solid phase reaction.

### Electric transport

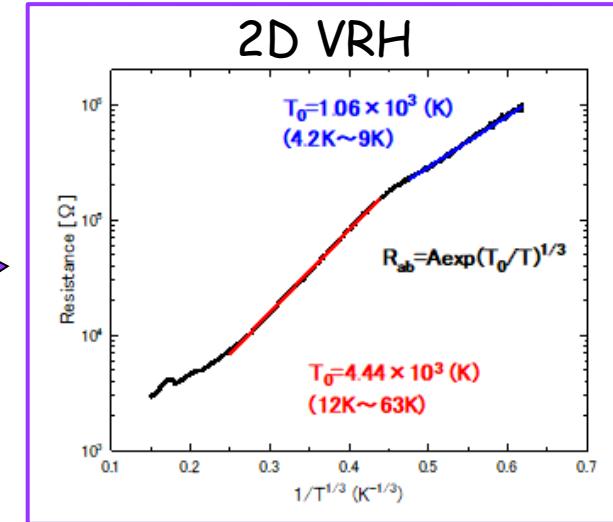
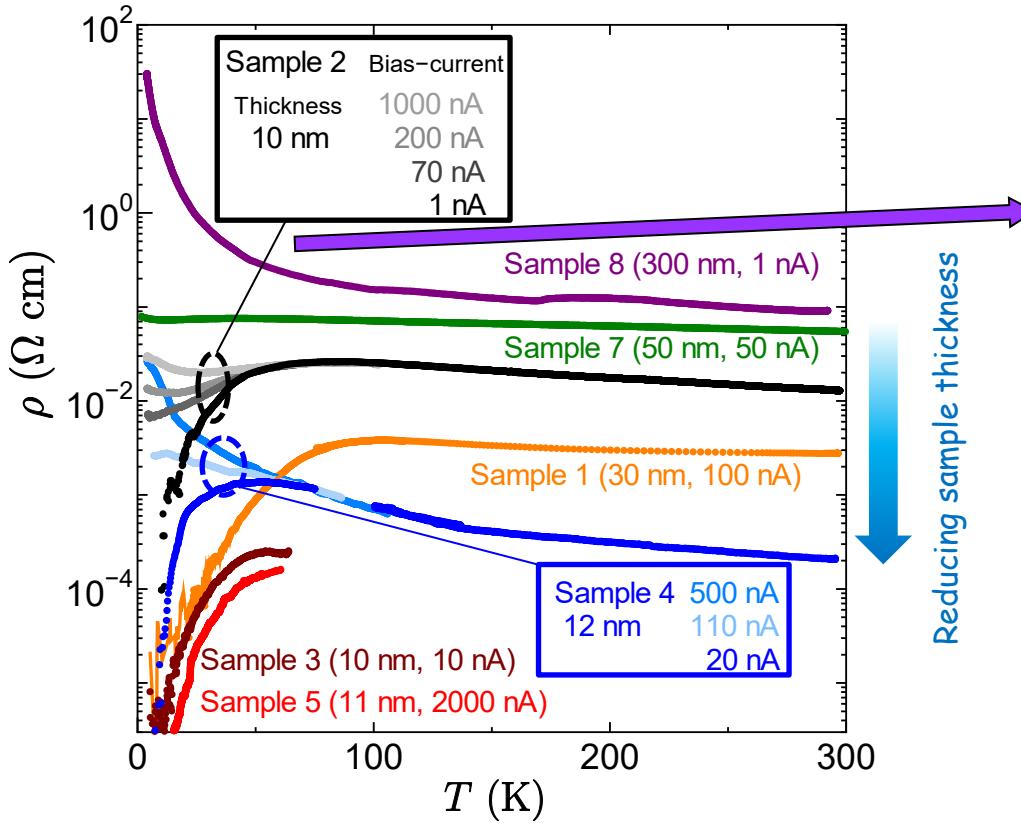
Attached gold electrodes  
by EB lithography



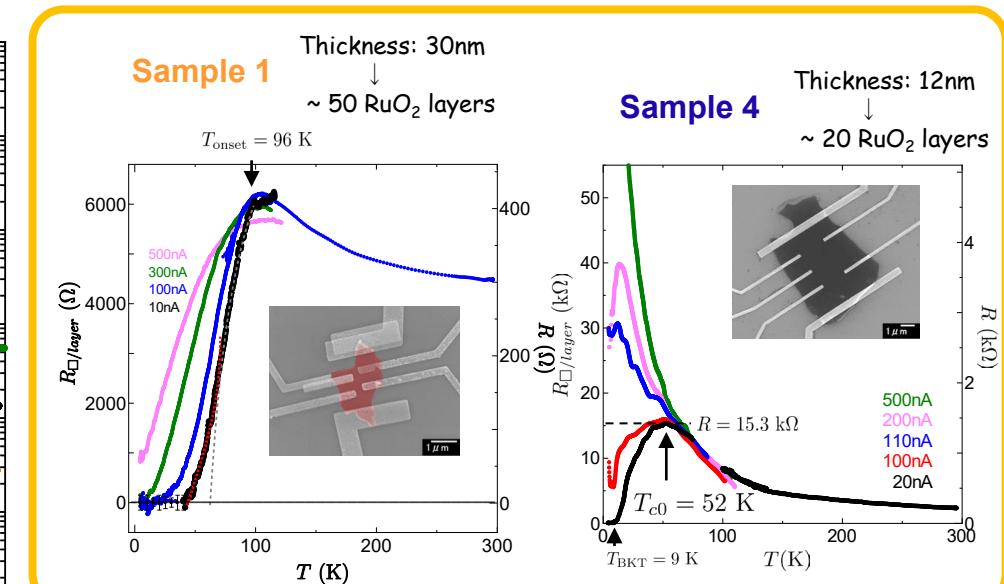
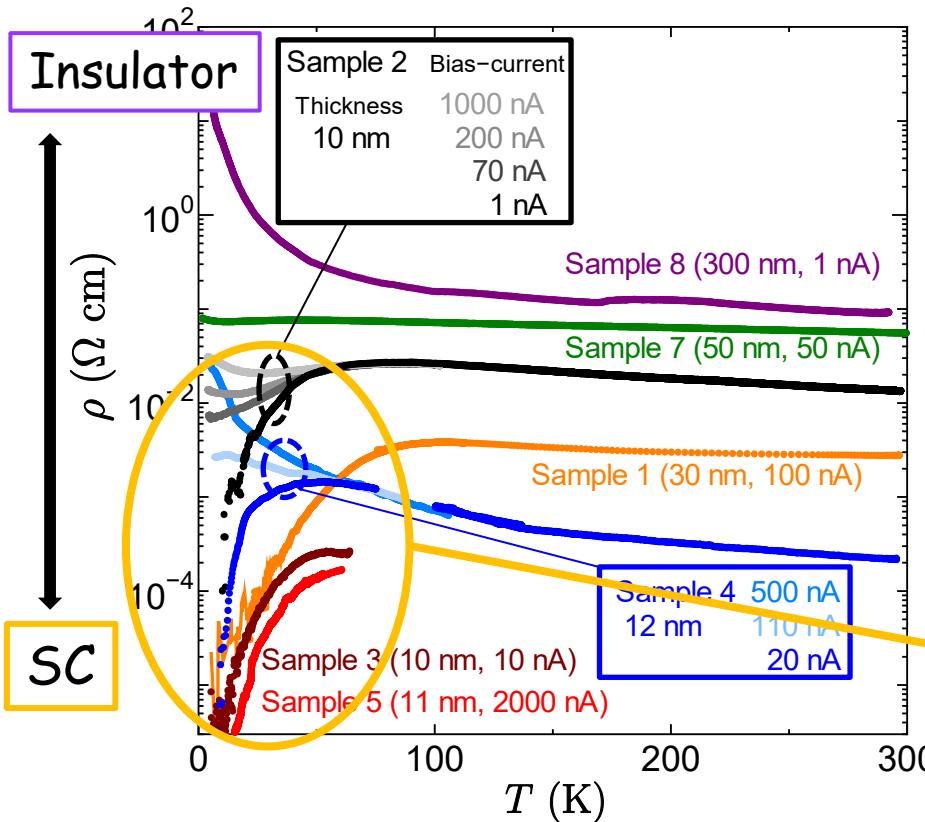
Powder consisting of nanoscale crystals

### Magnetization

# $\rho - T$ for different thicknesses



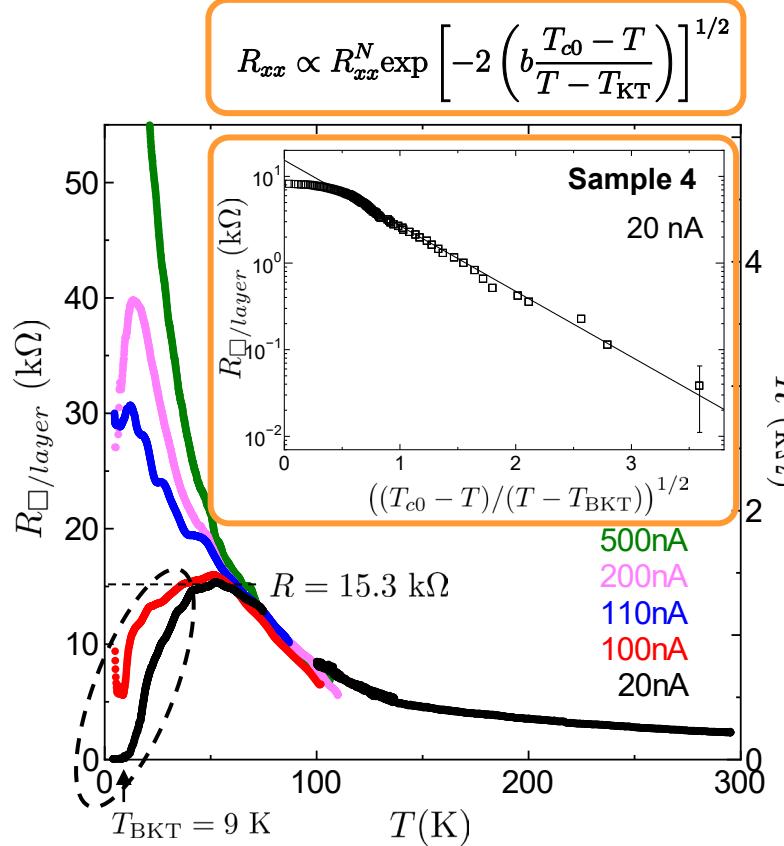
# Thickness-tuned SIT



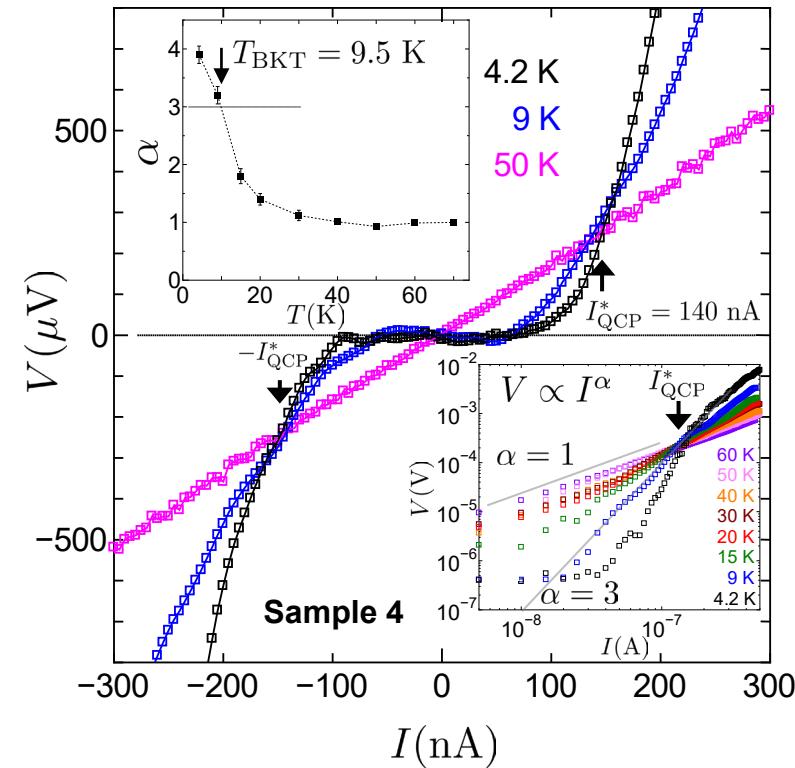
Resistivity drop  
Zero resistivity

We found the thickness-tuned SIT in nanoscale  $\text{Ca}_2\text{RuO}_4$ .

# BKT transition in nano- $\text{Ca}_2\text{RuO}_4$

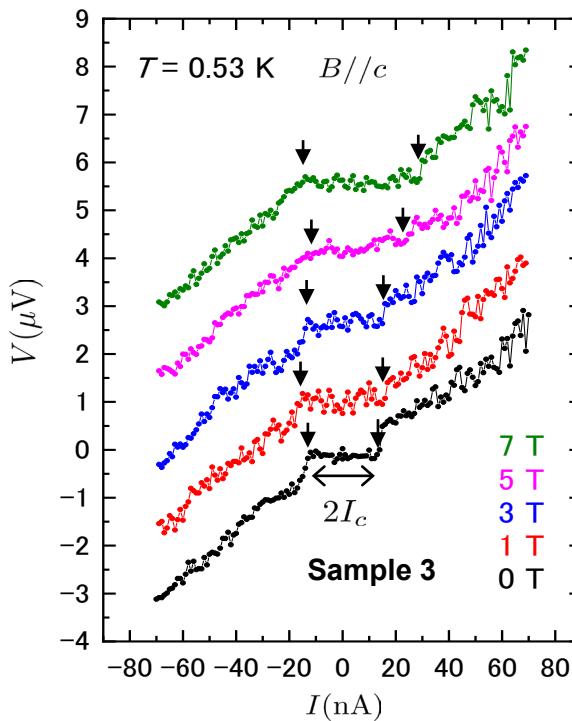


Broad transition due to vortex flow

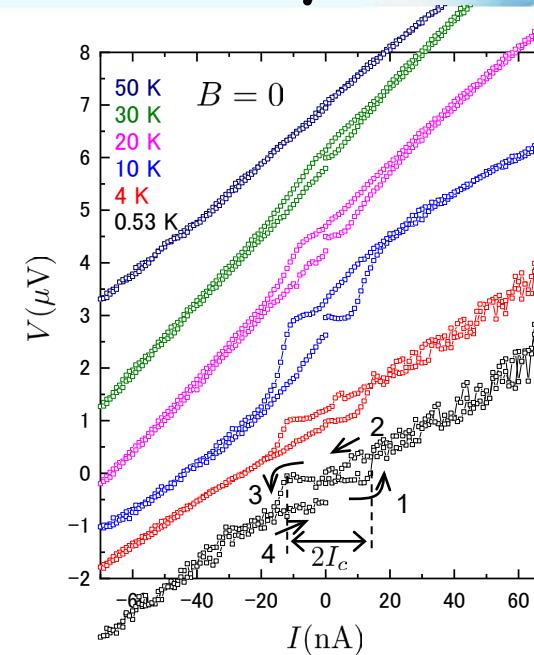


Universal jump of  $V \propto I^\alpha$

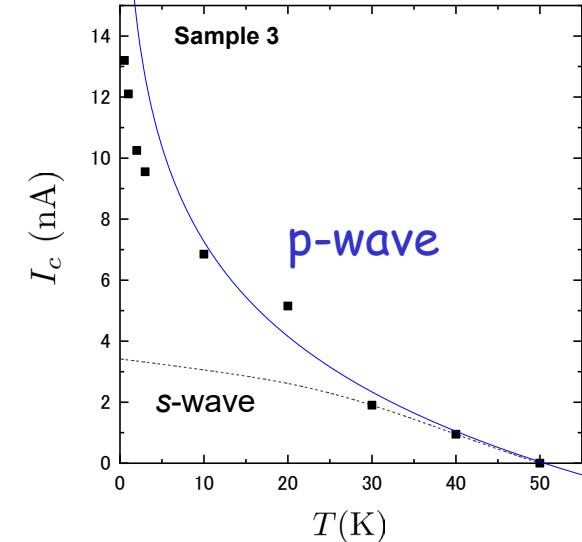
# Enhanced superconductivity



➤ Enhanced supercurrent



➤ Hysteresis behavior



➤  $I_c$  anomaly

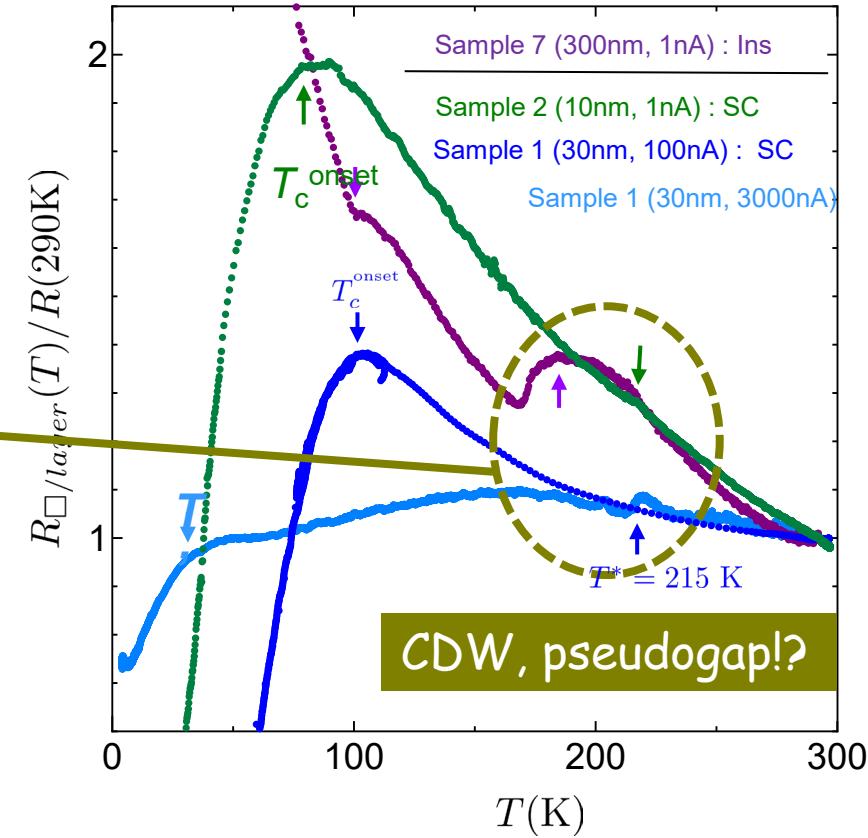
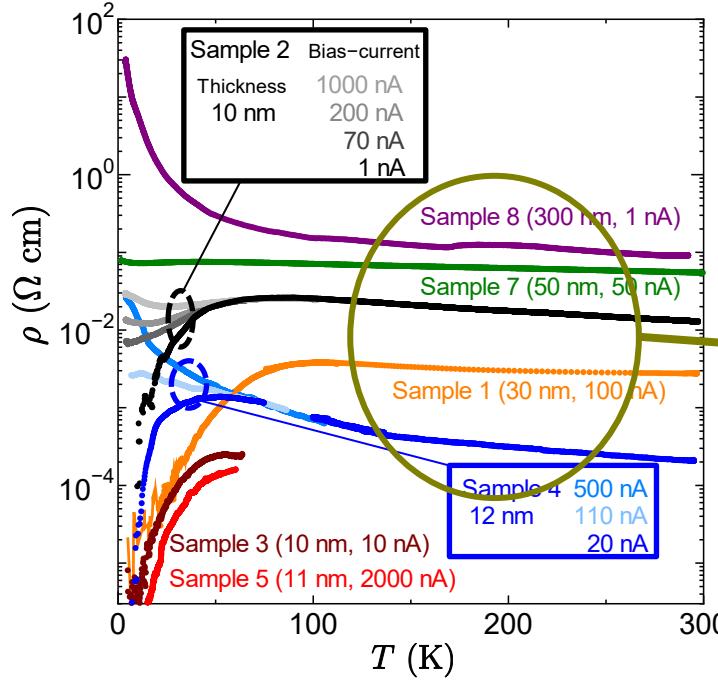
Josephson critical current  
between chiral p-wave SCs

$$I_c(T) = a I_{c0} \ln \left( \frac{b\Delta}{T} \right) \sin \Phi$$

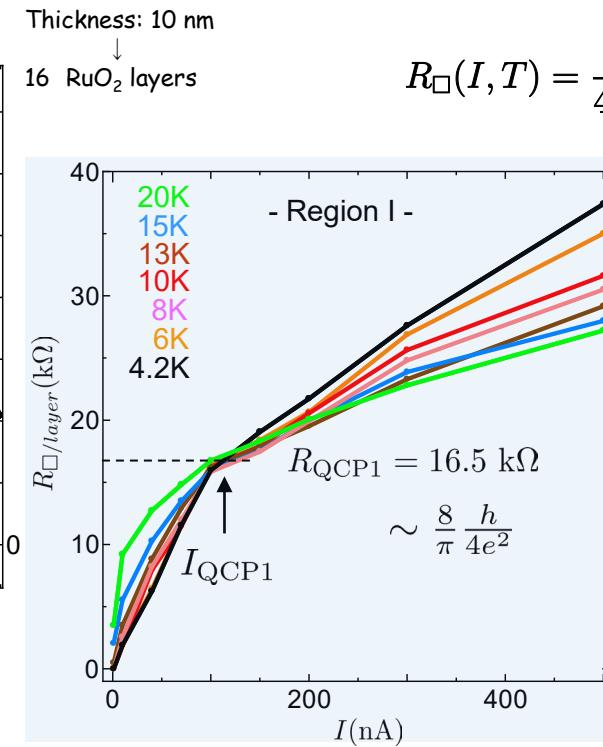
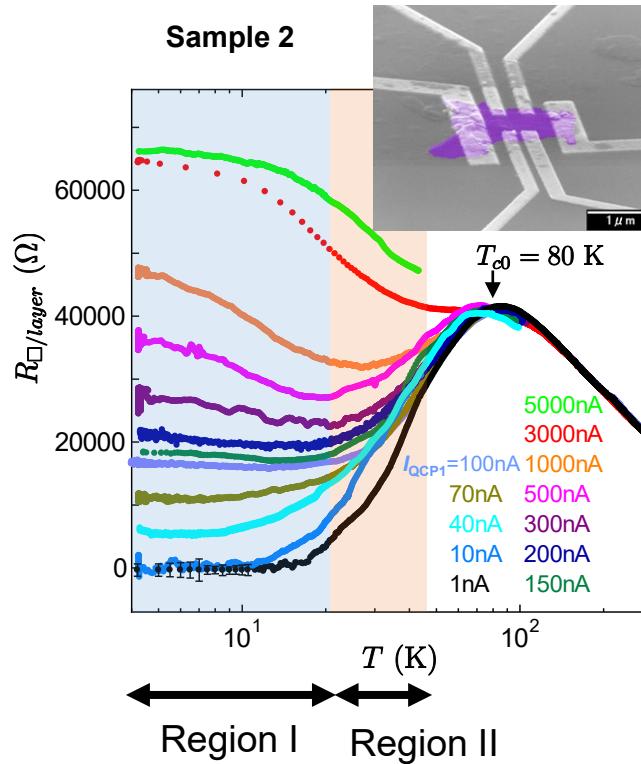
The superconductivity realizes the chiral p-wave state.

PSJ (2002)

# Resistance anomalies near 200 K

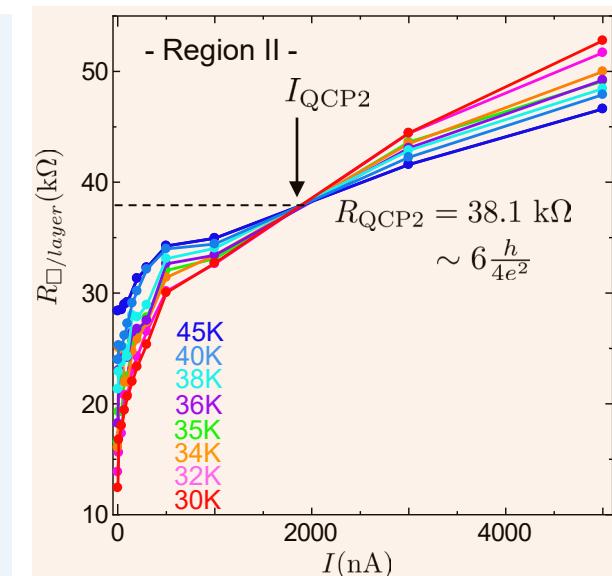


# Current-driven SIT

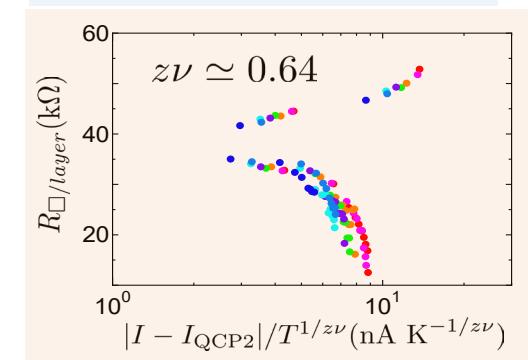
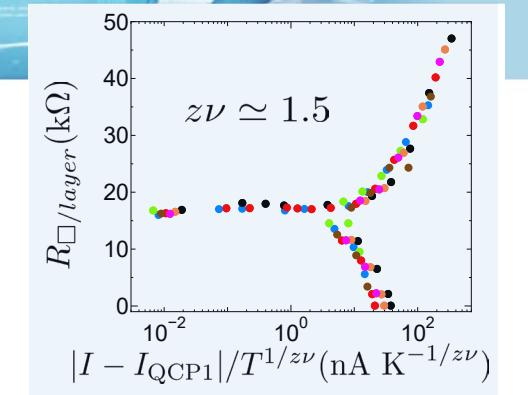
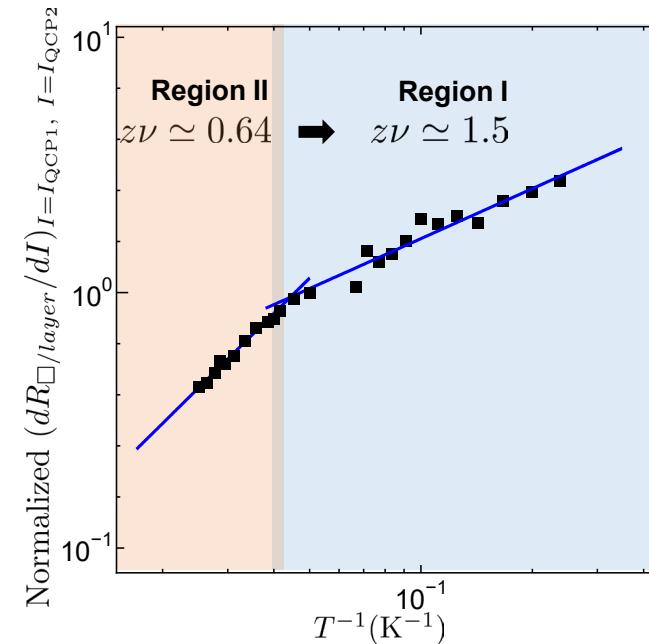
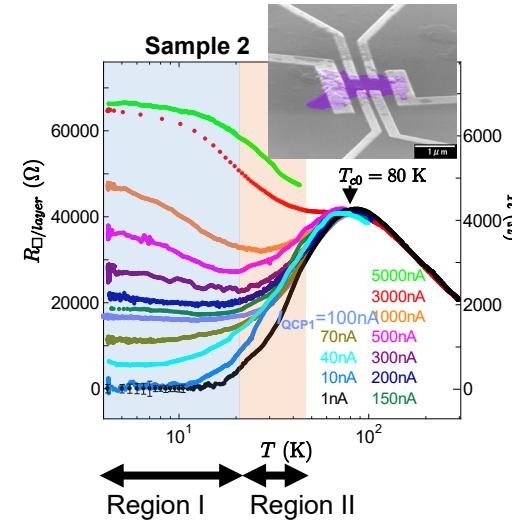


## Universal scaling analysis

$$R_{\square}(I, T) = \frac{h}{4e^2} f \left[ \frac{c_0(I - I_c)}{T^{1/z\nu}} \right]$$



# Two-stage quantum criticality

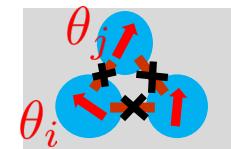


High  $T$      $z\nu = 2/3$     (2+1)D XY model in the clean limit



Low  $T$      $z\nu = 3/2$     Intrinsic inhomogeneities (quantum disorder)  
in strongly correlated systems

"Phase fluctuations"



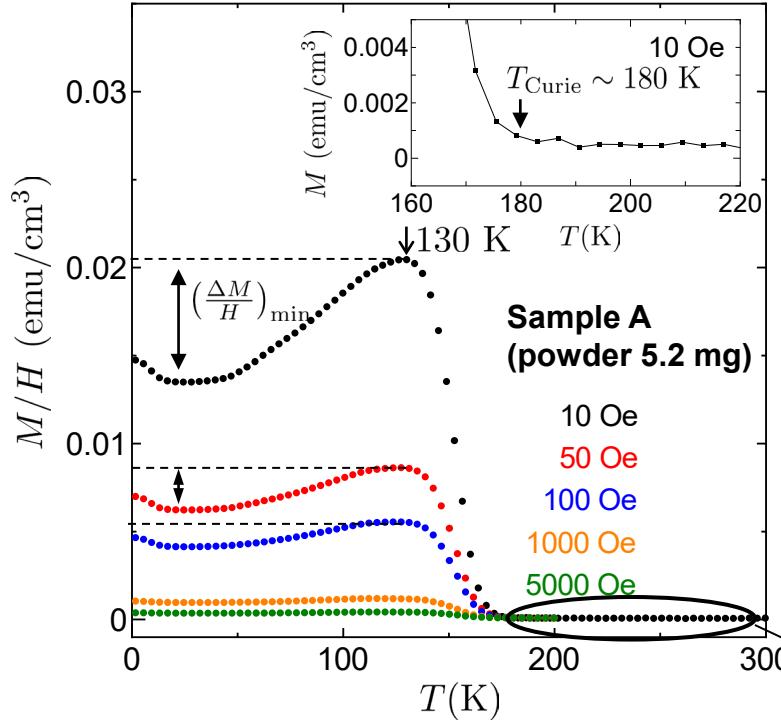
SC domains  
SC phase  
Josephson links

To clarify the diamagnetism,  
we performed magnetic measurements  
for powders consisting of nanoscale crystals.

Sample A : 5.2 mg  
Sample B : 2.6 mg

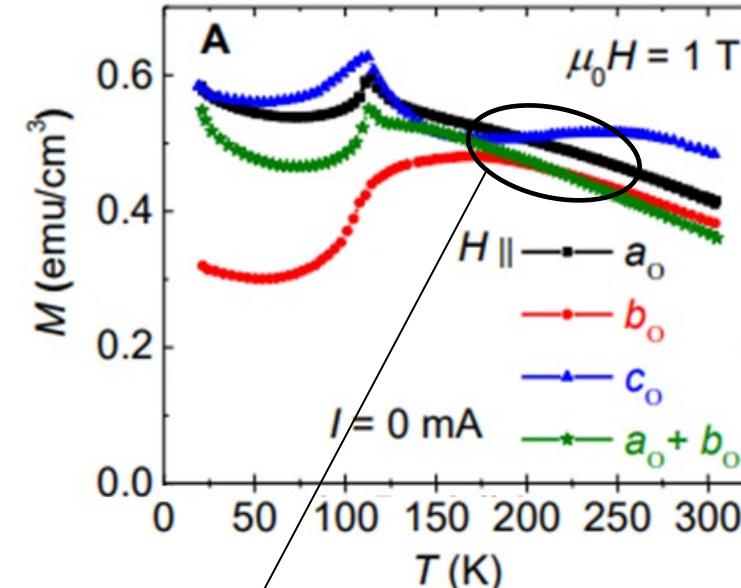
Temperature: 2K~300K  
Magnetic field :-7T~+7T  
Performed by MPMS3

# Magnetic properties



- $T_{\text{Curie}} = 180$  K
- Diamagnetic component

Bulk  $\text{Ca}_2\text{RuO}_4$  (AF insulator  $T_N = 113$ K)

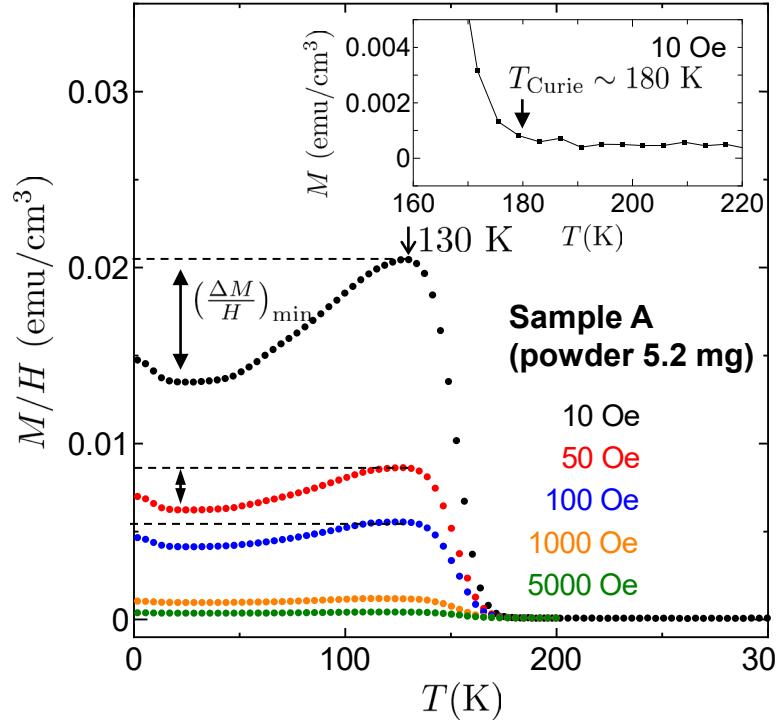


C. Sow *et al.*, Science. 358.1084(2017)

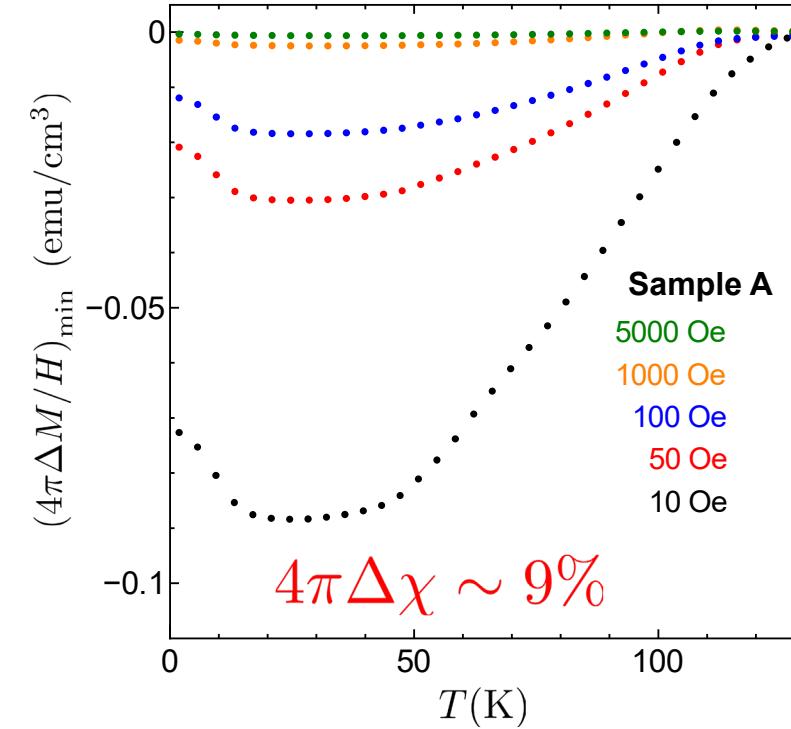
$$\chi \sim 5.0 \times 10^{-5} \text{ emu/cm}^3$$

Paramagnetic state

# Diamagnetism



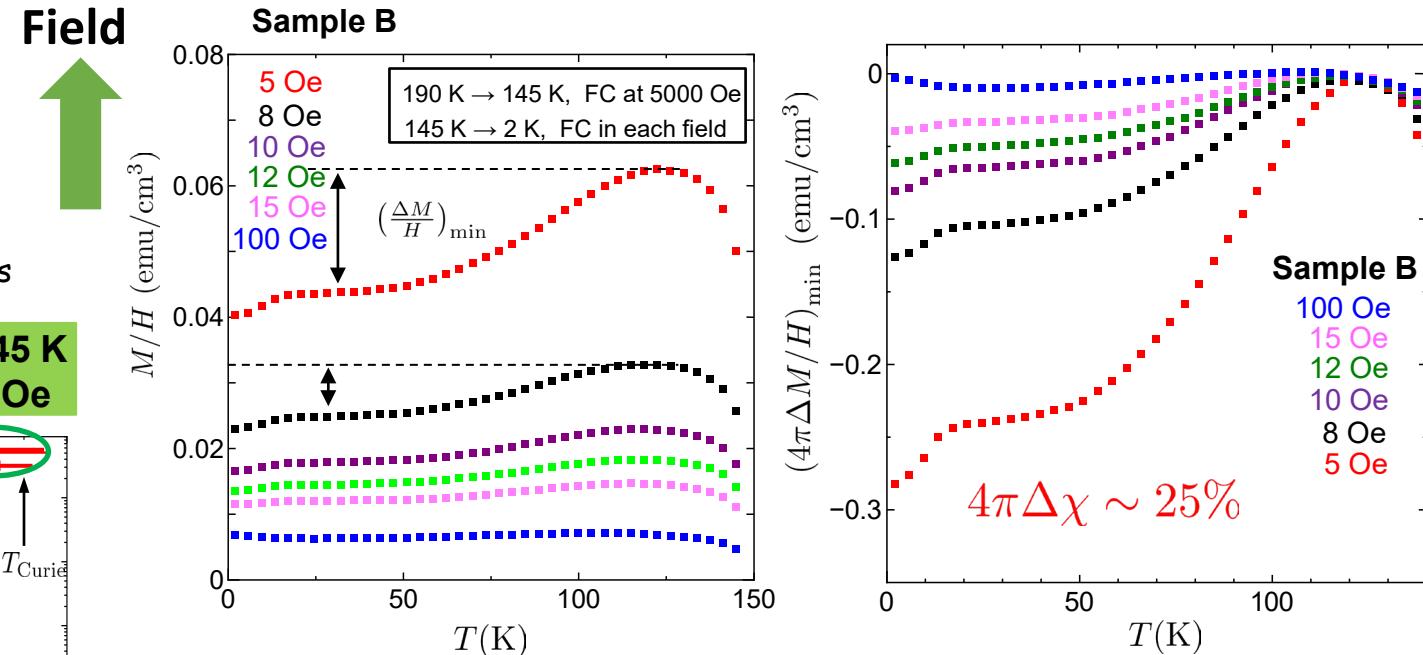
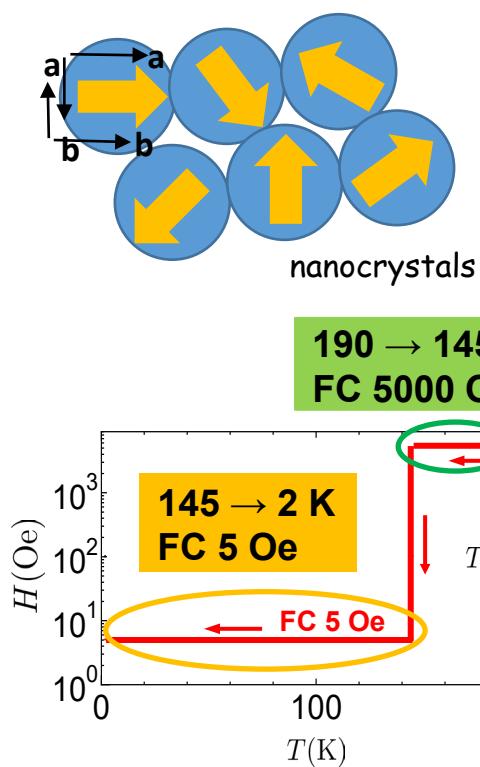
- $T_{\text{Curie}} = 180$  K
- Diamagnetic component



There are no phenomena that show such a giant diamagnetism other than superconductivity.

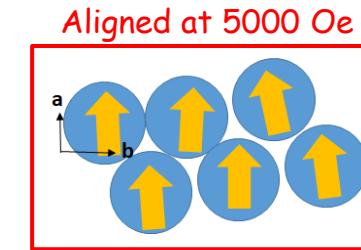
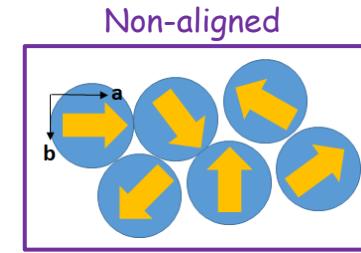
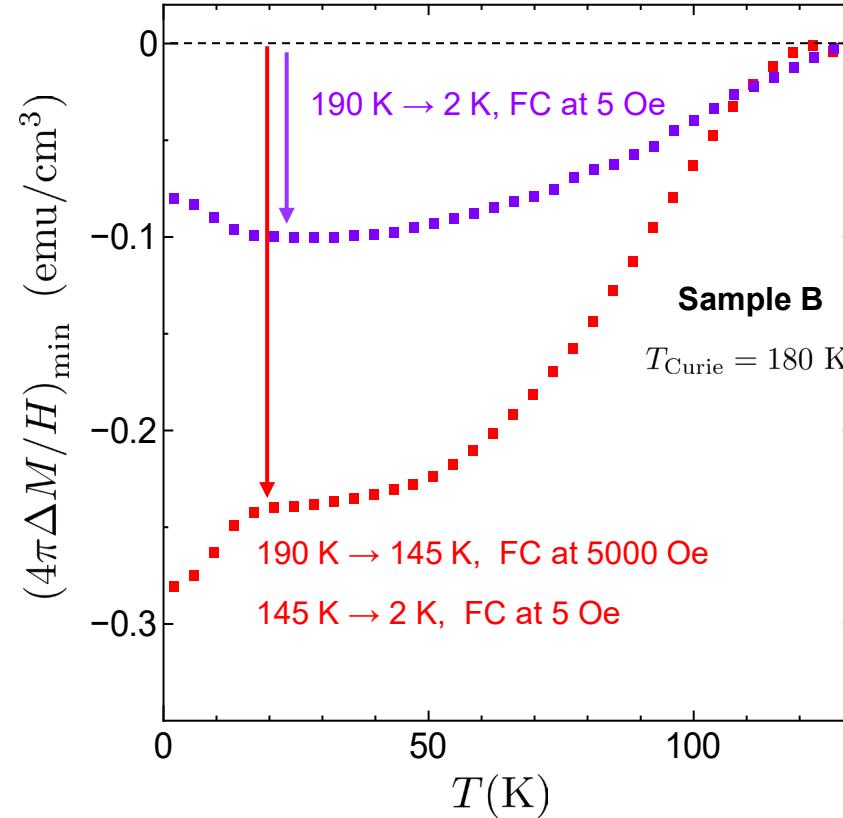
# Enhanced diamagnetic components

## Aligned magnetic domain



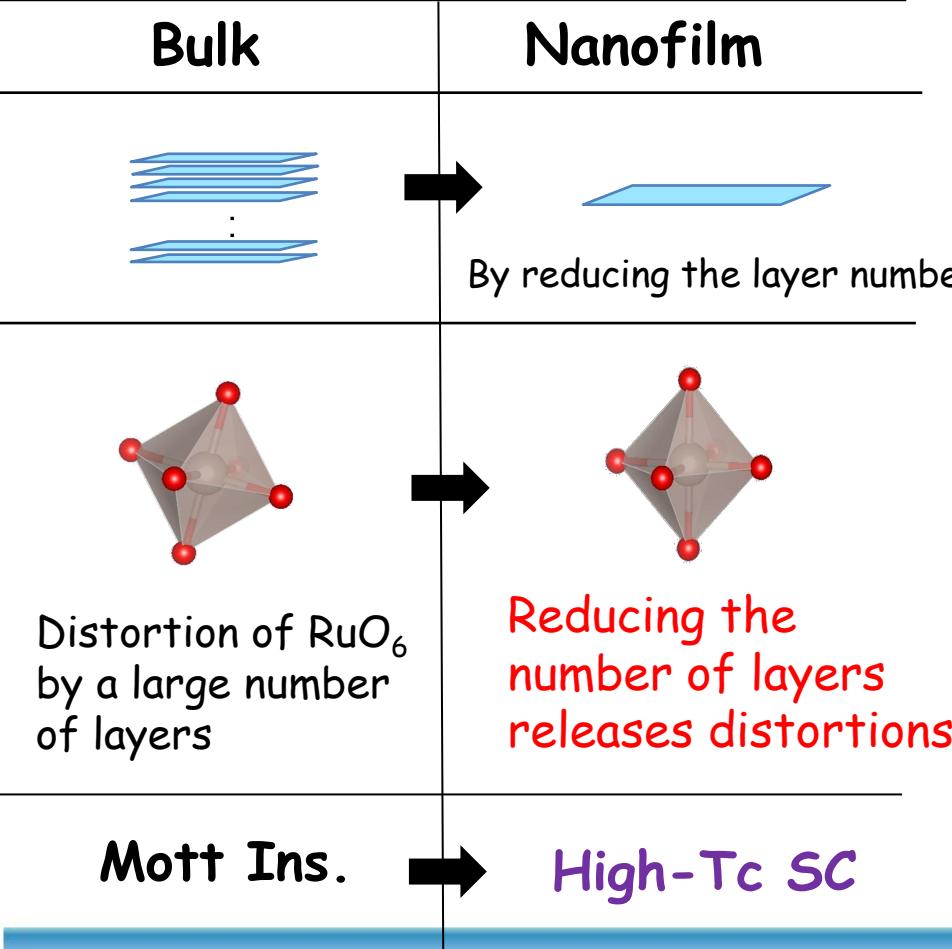
The diamagnetic components is  
enhanced by ferromagnetic ordering.

# Enhanced diamagnetism



Coexistence of superconductivity and ferromagnetism.

# Released distortion in nanocrystals

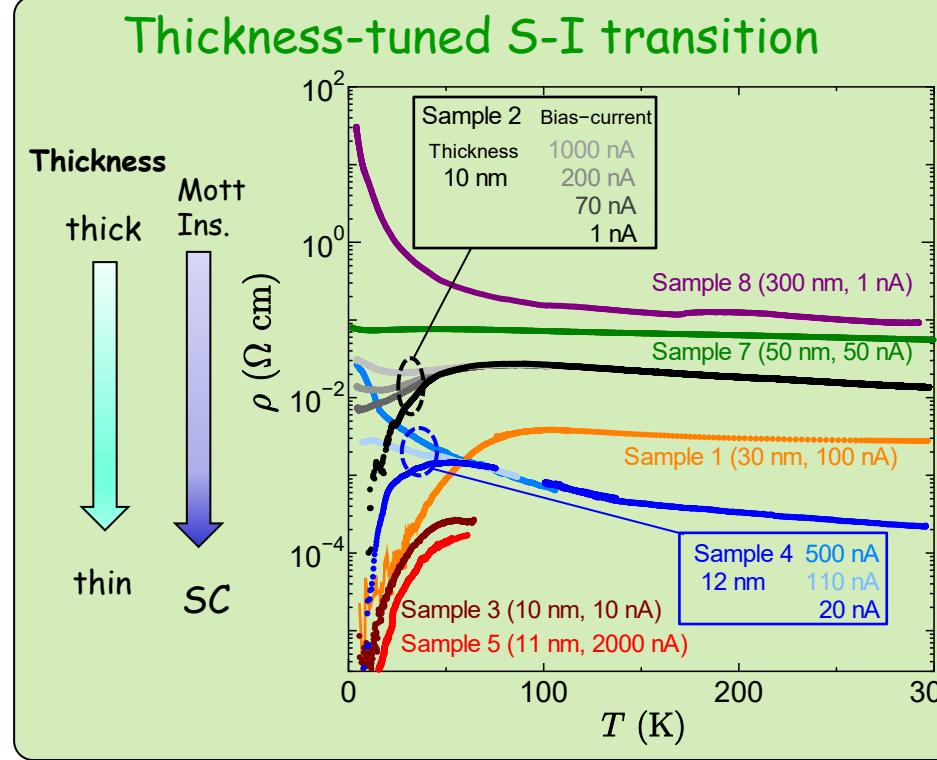


## First-principle calculation

	$a$ (Å)	$b$ (Å)	$c$ (Å)	Tilting (°)
Monolayer	5.806	4.847	—	10.01
Bi-layer	5.719	5.049	11.692	11.06
Tri-layer	5.694	5.128	11.65	11.34
Bulk	5.592	5.253	11.548	11.65

"Flattening" and "tilting" are released in nanofilms.

# Emergence of SC with negative pressure



## First-principle calculation

	$a$ (Å)	$b$ (Å)	$c$ (Å)	Tilting (°)
Mono-layer	5.806	4.847	—	10.01
Bi-layer	5.719	5.049	11.692	11.06
Tri-layer	5.694	5.128	11.65	11.34
Bulk	5.592	5.253	11.548	11.65

**"Flattening" and "tilting" are released in nanofilms.**

Suppression of the distortion under the negative pressure is responsible for SC behavior in CRO nanofilms.

# Summary

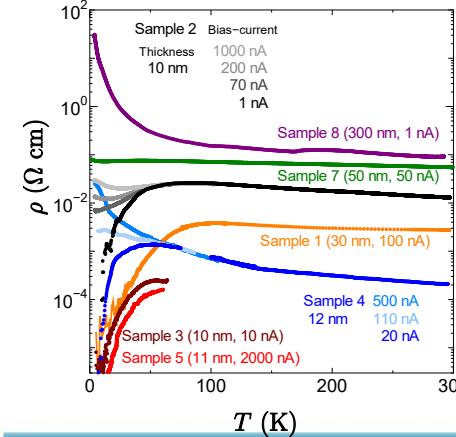


In 2D ruthenates, by reducing the number of layers, we found new ground states not observed in bulk.

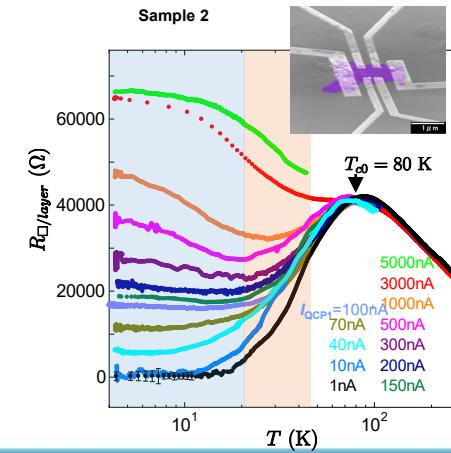
Coexistence of high- $T_c$  SC and ferromagnetism in nanoscale  $\text{Ca}_2\text{RuO}_4$

## Transport

◆ Thickness-tuned SIT

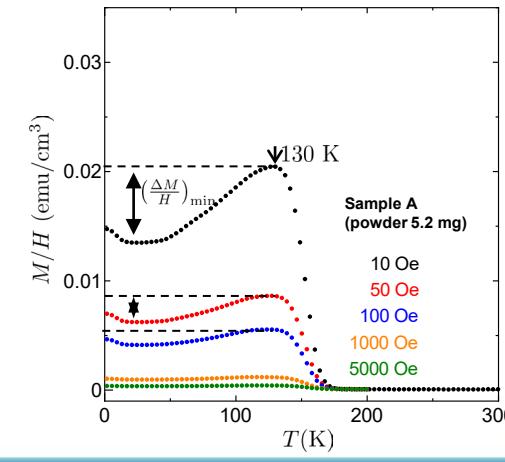


◆ Current-driven SIT



## Magnetism

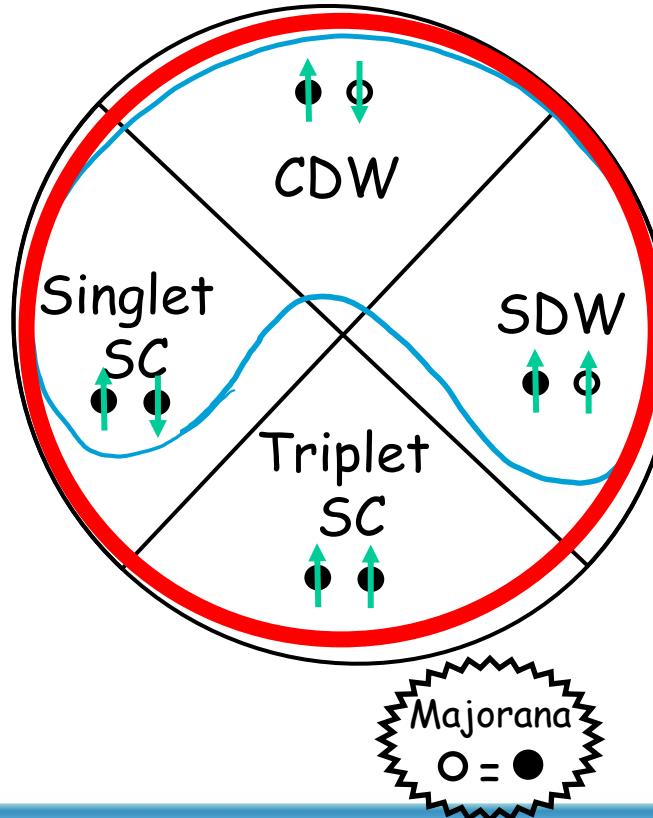
◆ SC coexisting with ferromagnetism



# Summary

- Electron
- Hole

## Ruthenates



Research for **ruthenate thin films** reveals that **new physics can be explored** distinct from in  $MX_2$  and cuprates.

