

CORRELATION FUNCTIONS FOR SOME CONFORMAL THEORIES ON RIEMANN SURFACES

MICHAEL MONASTYRSKY

*Institute of Theoretical and Experimental Physics,
117259 Moscow, Russia
e-mail: monastyrsky@vcitep.itep.ru*

SERGEI NECHAEV

*Landau Institute for Theoretical Physics,
117940, Moscow, Russia
e-mail: nechaev@landau.ac.ru*

Abstract

We discuss the geometrical connection between 2D conformal field theories, random walks on hyperbolic Riemann surfaces and knot theory. For the wide class of CFT's with monodromies being the discrete subgroups of $SL(2, \mathbb{R})$ the determination of four-point correlation functions are related to construction of topological invariants for random walks on multipunctured Riemann surfaces.

Introduction

The behavior of systems consisting of chain-like objects can be modified in an essential way by topological constraints. The topological problems have been widely investigated in connection with quantum field, string theories and quantum Hall effect [1], vortices in superfluid liquids (^3He , ^4He) [2] and thermodynamic properties of entangled polymers [3, 4] etc. In the last few years there has been much progress in understanding the relationship between Chern–Simons conformal field theory on the one hand and constructions of knot and link invariants on the other hand (see, for general references [5, 6]). However, despite the general concepts have been well elaborated in the field-theoretic context, their application to the related areas of mathematics and physics has been highly limited.

In this note we indicate how the simple geometrical constructions originating from the theory of representations of braid groups enable us to unite such at first sight different objects like CFT, random walks on hyperbolic surfaces and knot invariants. The main physical application due to this approach is the evaluation of four–point correlation functions in CFT on the Riemann surface associated with the discrete subgroups of the group $SL(2, \mathbb{R})$.

1 Random Walk on Punctured Plane and Conformal Field Theory

We start with the random walk of length L and the elementary step ℓ ($\ell \equiv 1$) on the complex plane $\mathcal{C}_1 = \{z | z = x + iy\}$ with two punctures. Suppose the coordinates of these points being M_1 ($z_1 = (0, 0)$) and M_2 ($z_2 = (a, 0)$) ($a \equiv 1$) (such choice does not indicate the loss of generality). Consider two closed paths on \mathcal{C}_1 and attribute the generators g_1, g_2 of some group G to the windings around the points M_1 and M_2 . if we move along the path in the clockwise direction (we apply g_1^{-1}, g_2^{-1}). The question is: what is the probability $P(\mu, L)$ for the random walk of length L on the plane \mathcal{C}_1 to form a closed loop with the primitive word written in terms of generators $\{g_1, g_2, g_1^{-1}, g_2^{-1}\}$ to have the length μ .

The function $P(\mu, L)$ can be written as a path integral with a Wiener measure

$$P(\mu, L) = \frac{1}{\Theta} \int \dots \int \mathcal{D}\{z\} \exp \left\{ - \int_0^L \left(\frac{dz(s)}{ds} \right)^2 ds \right\} \delta [W\{g_1, g_2, g_1^{-1}, g_2^{-1}|z\} - \mu] \quad (1)$$

where $\Theta = \int P(\mu, L) d\mu$ and $W\{\dots|z\}$ is the length of the primitive word on G as a functional of the path on the complex plane \mathcal{C}_1 . Conformal methods enable us to construct the topological invariant W for the group G .

To construct explicitly the topological invariants W for general group G we go to the universal covering of double punctured plane \mathcal{C}_1 . This covering is isomorphic to the upper half-plane $\mathcal{H} = \{\zeta | \zeta = \xi + i\lambda \lambda > 0\}$. The fundamental domain of the group $G\{g_1, g_2\}$ has the form of the curvilinear triangle in \mathcal{H} . Each fundamental domain represents the Riemann sheet corresponding to the fibre bundle above \mathcal{C}_1 . The universal covering space $\mathcal{Z}(\zeta)$ is the unit of all such Riemann sheets.

Coordinates of ends of the trajectory on \mathcal{H} can be served as the topological invariant for the path on double punctured plane \mathcal{C}_1 with respect to the action of the group G . It should be emphasized that this invariant is *complete* for our purposes. In particular, coordinates of initial and final points of any trajectory on \mathcal{H} determine: (a) the coordinates of corresponding points on \mathcal{C}_1 ; (b) the homotopy class of any path on \mathcal{C}_1 : The paths on \mathcal{H} are closed if and only if $W\{g_1, g_2|z\} \equiv 1$, i.e. they belong to the trivial homotopy class. After Fourier transform $P(\mu, L) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-ip\mu} P(p, L) dp$ the function $P(p, L)$ coincides with the Green function $P(\mathbf{r}_0, \mathbf{r} = \mathbf{r}_0, p, L)$ of the non-stationary Schrödinger-like equation

$$\frac{\partial}{\partial L} P(\mathbf{r}_0, \mathbf{r}, p, L) - \left(\frac{1}{2} \nabla - ip\mathbf{A}(\mathbf{r}) \right)^2 P(\mathbf{r}_0, \mathbf{r}, p, L) = \delta(L) \delta(\mathbf{r} - \mathbf{r}_0) \quad (2)$$

for the free particle motion in a "magnetic field" with the vector potential $\mathbf{A}(\mathbf{r})$ defined by the monodromy properties of the group G ; p plays the role of a "charge" and the magnetic field is transversal, i.e. $\text{rot}\mathbf{A}(\mathbf{r}) = 0$. On the complex plane \mathcal{C}_1 the field $A(z)$ is determined as $A(z) = d\zeta(z)/dz$.

Now let us construct the desired conformal mapping. The single-valued function $z(\zeta) \equiv \zeta^{-1}(z)$ is defined in the fundamental domain of the group G on \mathcal{H} . The multivalued function $\phi(\zeta)$ is determined as follows: (a) in the fundamental domain $\phi(\zeta) = z(\zeta)$; (b) in all other domains on \mathcal{H} the function $\phi(\zeta)$ is analytically continued through the boundaries by means of the fractional transformations that are generated by G .

Consider two basic contours P_1 and P_2 on \mathcal{H} being the conformal images of the contours $A_1(z)$ and $A_2(z)$ corresponding to the generators g_1 and g_2 of G . The function $\phi(z)$ ($z \neq \{z_1, z_2, \infty\}$) obeys the following transformations:

$$\phi \left[z \xrightarrow{A_1} z \right] \rightarrow \tilde{\phi}_1(z) = \frac{a_1\phi(z) + b_1}{c_1\phi(z) + d_1}; \quad \phi \left[z \xrightarrow{A_2} z \right] \rightarrow \tilde{\phi}_2(z) = \frac{a_2\phi(z) + b_2}{c_2\phi(z) + d_2} \quad (3)$$

where

$$\begin{pmatrix} a_1 & b_1 \\ c_1 & d_1 \end{pmatrix} = \hat{g}_1; \quad \begin{pmatrix} a_2 & b_2 \\ c_2 & d_2 \end{pmatrix} = \hat{g}_2 \quad (4)$$

are the matrices of basic substitutions of the group $G\{g_1, g_2\}$.

Assuming $\zeta(z)$ to be a ratio of two fundamental solutions, $u_1(z)$, and, $u_2(z)$, of some second order differential equation with branching points $\{z_1 = 0, z_2 = 1, z_3 = \infty\}$, we conclude that the solutions $u_1(z)$ and $u_2(z)$ undergo the linear transformations when the variable z moves along the contours $A_1(z)$ and $A_2(z)$:

$$A_1(z) : \begin{pmatrix} \tilde{u}_1(z) \\ \tilde{u}_2(z) \end{pmatrix} = \hat{g}_1 \begin{pmatrix} u_1(z) \\ u_2(z) \end{pmatrix}; \quad A_2(z) : \begin{pmatrix} \tilde{u}_1(z) \\ \tilde{u}_2(z) \end{pmatrix} = \hat{g}_2 \begin{pmatrix} u_1(z) \\ u_2(z) \end{pmatrix} \quad (5)$$

The problem of restoring the form of differential equation from the monodromy matrices g_1 and g_2 of the group G is known as Riemann–Hilbert problem and has a long history [8, 9].

Let us consider the special class of triangle groups—so-called Hecke groups $H(h)$:

$$T(h) = \begin{pmatrix} 1 & h \\ 0 & 1 \end{pmatrix}; \quad S = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \quad (6)$$

where $h = 2 \cos \frac{\pi}{q}$ and $q = 3, 4, 5, 6, \dots$ (the case $h = 1$ corresponds to $PSL(2, \mathbf{Z})$). The fundamental domain of the Hecke group is the circular triangle ABC with angles $\{0, \frac{\pi}{q}, \frac{\pi}{2}\}$ lying on \mathcal{H} .

The function $\zeta(z, q) = u_1(z, q)/u_2(z, q)$ performs the conformal mapping of $\text{Im}z > 0$ with three punctures $(0, 1, \infty)$ to the interior of the triangle ABC on \mathcal{H} , where $u_1(z, q)$ and $u_2(z, q)$ are the fundamental solutions of Picard-Fuchs hypergeometric equation

$$z(z-1)u''(z, q) + \left[\left(\frac{3}{2} - \frac{1}{q} \right) z - \left(1 - \frac{1}{q} \right) \right] u'(z, q) + \frac{1}{4} \left(\frac{1}{2} - \frac{1}{q} \right)^2 u(z, q) = 0 \quad (7)$$

The former geometrical construction is evidently related to CFT. Eq.(7) can be associated with the equation on the four-point correlation function of some critical CFT. The question remains *whether it is always possible to adjust the parameters of the corresponding critical CFT to the coefficients of equations like (7)*. In some particular cases this question has positive answer and we describe briefly the general construction.

Consider CFT characterized by the parameters $\{\Delta_1, \Delta_2, \Delta_3 = \Delta_1 + \Delta_2 - \Delta_{(12)}, \delta\}$ which determine the conformal dimensions [11]. Suppose that $\varphi(z)$ is the primary field. The four-point correlation function $\langle \varphi(z_1)\varphi(z_2)\varphi(z_3)\varphi(z) \rangle$ satisfies the partial differential equation, that follows from the conformal Ward identity [11, 12]. Introducing anharmonic ratio in the standard way, we get

$$\left[\frac{3}{2(2\delta + 1)} \frac{d^2}{dz^2} + \left(\frac{1}{z} + \frac{1}{z-1} \right) \frac{d}{dz} - \left(\frac{\Delta_1}{z^2} + \frac{\Delta_2}{(z-1)^2} - \frac{\delta + \Delta_{(12)}}{z(z-1)} \right) \right] \psi(z|0, 1, \infty) = 0 \quad (8)$$

Recall, that $\delta \equiv \Delta_{(n,m)} = \frac{1}{24}(c-1) + (\frac{1}{2}\alpha_+n + \frac{1}{2}\alpha_-m)^2$ where $\alpha_{\pm} = \frac{\sqrt{1-c} \pm \sqrt{25-c}}{\sqrt{24}}$; (n, m) are positive integers and c is the central charge of corresponding Virasoro algebra.

The substitution $\psi(z|0, 1, \infty) = z^{-2b_-} (z-1)^{-2b_+} u(z)$ where $b_{\pm} = \frac{1}{6}(2\delta + 1) - \frac{q \pm 1}{4q}$ allows us to rewrite (8) in the form of (7) by setting

$$\begin{cases} \Delta_1 = \frac{9 - q^2 + 8\delta q^2 - 16\delta^2 q^2}{24(1 + 2\delta)q^2} \\ \Delta_2 = \frac{9 - 18q + 5q^2 + 32\delta q^2 - 64\delta^2 q^2}{96(1 + 2\delta)q^2} \\ \Delta_{(12)} = \frac{1}{2\delta + 1} \left[\frac{9}{8q^2} - \frac{9}{8q} + \frac{\delta}{2} \left(1 - \frac{1}{q} \right) - \frac{10}{3}\delta^2 + \frac{1}{3} \right] \end{cases} \quad (9)$$

For $q = 3$ we always have $\Delta_1 = \Delta_2$ and one of possible solutions (corresponding to $c = -2$) gives us $\delta = -\frac{1}{8}$, $\Delta_1 = \Delta_2 = -\frac{5}{72}$, $\Delta_{(12)} = -\frac{1}{72}$.

2 Structure of Target Space of Braid Group B_3 and Random Walk on B_3

The braid group B_n of n strings has $n-1$ generators $\{\sigma_1, \sigma_2, \dots, \sigma_{n-1}\}$ with the following relations:

$$\begin{aligned}\sigma_i \sigma_{i+1} \sigma_i &= \sigma_{i+1} \sigma_i \sigma_{i+1} & (1 \leq i < n-1) \\ \sigma_i \sigma_j &= \sigma_j \sigma_i & (|i-j| \geq 2) \\ \sigma_i \sigma_i^{-1} &= \sigma_i^{-1} \sigma_i = e\end{aligned}\tag{10}$$

Let us mention that the *length* of the braid is the total number of used letters, while the *minimal irreducible length* hereafter referred to as the "primitive word" is the shortest noncontractible length of a particular braid which remains after applying all possible group relations (10). Any braid corresponds to some knot or link. Hence, it is feasible to use the braid group representation for the construction of topological invariants of knots and links. However the correspondence between braids and knots is not one-to-one and each knot or link can be represented by infinite numbers of different braids.

The (right-hand) random walk (the random word) on the group B_n with a uniform distribution over generators $\{\sigma_1, \dots, \sigma_{n-1}, \sigma_1^{-1}, \dots, \sigma_{n-1}^{-1}\}$ means that with the probability $\frac{1}{2n-2}$ we add the element σ_{α_N} or $\sigma_{\alpha_N}^{-1}$ to the given word of $N-1$ generators (letters) from the right-hand side. The most attention should be paid to the question: what is the probability $P(\mu, N)$ to find the μ -letter primitive word in the N -letter random word in the group B_n . In [10] it is conjectured that $P(\mu, N)$ obeys the following asymptotics

$$P(\mu, N) \propto \frac{\mu}{N^{3/2}} \exp\left(-a(n)N + b(n)\mu - \frac{\mu^2}{c(n)N}\right)\tag{11}$$

where $a(n)$, $b(n)$, $c(n)$ are numerical constants depending on n only.

The group B_3 can be represented by $PGL(2, \mathbf{R})$ -matrices. To be more specific, the braid generators σ_1 and σ_2 in the Burau representation [7] read

$$\sigma_1 = \begin{pmatrix} -t & 1 \\ 0 & 1 \end{pmatrix}; \quad \sigma_2 = \begin{pmatrix} 1 & 0 \\ t & -t \end{pmatrix}\tag{12}$$

where t is the "spectral parameter". For $t = -1$ the matrices σ_1 and σ_2 generate $PSL(2, \mathbf{Z})$ such that B_3 is its central extension.

We define $\tilde{\sigma}_1 = \sigma_1$, $\tilde{\sigma}_2 = \sigma_2$ (at $t = -1$) and rewrite the generators of the modular group $PSL(2, \mathbf{Z})$ in the canonical form using the matrices S, T (compare to (6)).

Figure 1: The graph corresponding to the structure of: (a) the modular group $PSL(2, \mathbf{Z})$ (i.e. the Hecke group with $q = 3$) and (b) the Hecke group with $q = 4$.

The braiding relation $\tilde{\sigma}_1 \tilde{\sigma}_2 \tilde{\sigma}_1 = \tilde{\sigma}_2 \tilde{\sigma}_1 \tilde{\sigma}_2$ in the $\{S, T\}$ -representation takes the form $S^2 T S^{-2} T^{-1} = 1$ (we also have $S^4 = (ST)^3 = 1$).

The group $PSL(2, \mathbf{Z})$ is a discrete subgroup of $PSL(2, \mathbf{R})$. The fundamental domain of $PSL(2, \mathbf{Z})$ is a circular triangle with angles $\{0, \frac{\pi}{3}, \frac{\pi}{2}\}$ located in \mathcal{H} . The group $PSL(2, \mathbf{Z})$ is completely defined by its basic substitutions: (i) $S : \zeta \rightarrow -1/\zeta$; (ii) $T : \zeta \rightarrow \zeta + 1$. Let us choose an arbitrary element ζ_0 from the fundamental domain and construct a graph, $C(\Gamma)$ connecting the neighboring images of the initial element ζ_0 obtained under successive action of the generators from the set $\{S, T, S^{-1}, T^{-1}\}$. The graph $C(\Gamma)$ is shown in the Fig.1a. It will be seen that although this graph has the local cycles, its "backbone" has a Cayley tree ("free group") structure. Thus, the problem of limit distribution of a random walks on the modular group can be related to the problem of random walks on the graph $C(\Gamma)$. Hence, we immediately can compute the probability $P(\mu, N)$ that the N -step random walk along the graph $C(\Gamma)$ which starts from the origin ends in some vertex of the cell at the distance of μ steps along the backbone. This distribution has the form of Eq.(11).

Now we can construct the "target space" of the braid group B_3 . Introduce the normalized generators of the group B_3 : $\|\sigma_1\| = \tau(t)$; $\|\sigma_2\| = \tau^{-1}(t)s(t)\tau^{-1}(t)$, where

$\|\sigma_1(t)\| = (\det \sigma_1)^{-1} \sigma_1$ and $\|\sigma_2(t)\| = (\det \sigma_2)^{-1} \sigma_2$. By $\tau(t)$ and $s(t)$ we denote the generators of the "t-deformed" group $PSL_t(2, \mathbf{Z})$. The group $PSL_t(2, \mathbf{Z})$ keeps invariable the relations of $PSL(2, \mathbf{Z})$, namely, we have $(\tau(t)s(t))^3 = s^4(t) = \tau(t)s^2(t)\tau^{-1}(t)s^{-2}(t) = 1$. Hence, we ultimately come to the conclusion that the graphs $C(\Gamma_t)$ (for $PSL_t(2, \mathbf{Z})$) and $C(\Gamma)$ are topologically equivalent. However the metric properties of $C(\Gamma_t)$ and $C(\Gamma)$ are different because of different embeddings of $PSL_t(2, \mathbf{Z})$ and $PSL(2, \mathbf{Z})$ into the upper half-plane. The additional benefit of this construction is that we can represent the graph of B_3 on the graph $C(\Gamma)$ if we supply in each center of the cell the "magnetic pole" with the quantity of the flux equal to t^6 . Fig.1b represents the graph corresponding to Hecke group with $q = 4$.

To combine the monodromy representation of the Hecke group $H(h)$ with the topological properties of knots on the basis of braid group representation, consider the space of functions $V = \{f_{ij}\}$ where

$$f_{ij}(x_1, x_2, \dots, x_n) = \int_{x_i}^{x_j} (z - x_1)^{-\mu_1} \dots (z - x_n)^{-\mu_n} dz \quad (13)$$

μ_1, \dots, μ_n are the complex numbers and path integration is selected in the space $CP^1 \setminus \{x_1, \dots, x_n\}$. The group B_n acts on f_{ij} if $\mu_i = \mu_j$ but the group of colored braids C_n acts on the general linear span of the multivalued functions $\{f_{ij} : 0 \leq i < j \leq n - 1\}$.

In particular case of $n = 3$ we study the integral

$$\int \omega = \int z^{-\mu_1} (z - 1)^{-\mu_2} (z - x)^{-\mu_3} dz$$

Due to the classical result of H.A. Schwarz [13] we deduce that for any set of points $x_1, x_2 \notin \{0, 1, \infty, x\}$ the integral $\int_{x_1}^{x_2} \omega$ is the multivalued function of x satisfying the general hypergeometric equation

$$z(z - 1)u''(z) + \{(\alpha + \beta + 1)z - \gamma\}u'(z) + \alpha\beta u(z) = 0$$

where α, β, γ are arbitrary real constants related to μ_i . The analysis of this equation shows that the set of all monodromies arising by moving one puncture (x , for instance) around other ones ($0, 1, \infty$) generates the monodromy group $C_4(CP^1)$ on V . This group is discrete for some special constraint on exponents $\{\mu_i\}$ only (see [13]). Denote these monodromies by $A(x|0), A(x|1), A(x|\infty) = [A(x|0)A(x|1)]^{-1}$ (compare to (5)).

Let us briefly discuss the procedure of constructing the invariants of knots generated by the representations of C_4 . In the Jones' theory of polynomial invariants it has been

pointed out the remarkable connection between representations of II_1 -factors and knot invariants [15]. In particular it was proved that if the index value of factor II_1 is less than 4, there exists an integer $q \geq 3$ with index value $r = 4 \cos^2 \frac{\pi}{q}$. The II_1 -factor is hyperfinite.

To each element of the group $C_4(CP^1)$ we associate the word W represented by the monodromy matrices \hat{g}_1 and \hat{g}_2 . Any braid (a) generates some knot and (b) completely determines the monodromy matrix $M^{(N)} = \prod_{i=1}^N g_{\alpha_i}$, where $g_{\alpha_i} = \{g_1, g_2, g_1^{-1}, g_2^{-1}\}$ ($i = 1, 2, \dots, N$) and N is the total number generators used in the word W . Thus, tracing of $M^{(N)}$ gives the invariant of a knot. Using the classification of discrete subgroups of the $PSL(2, \mathbf{R})$ with two generators [14] (which includes also the Hecke and other triangle groups) we can expand our construction to the general case.

Concluding Remarks

1. Comparing the target spaces of the modular ($PSL(2, \mathbf{Z})$) and free (Γ_2) groups we see that the branching f_{eff} of the graph (i.e. the effective "curvature" of the corresponding hyperbolic space, in which the graph could be embedded) varies for different q . In particular, f_{eff} is changing from $f_{\text{eff}} = 3$ for $PSL(2, \mathbf{Z})$ to $f_{\text{eff}} = 4$ for Γ_2 .

2. Let us stress the that the index value r of II_1 -type factors, is equal to h^2 where h is the Hecke group parameter. Thus, the braids connected with the representations of II_1 -type factors can be realized as monodromies of Hecke groups and corresponding knot invariants could be produced by tracing of monodromy matrices.

3. Topological properties of knots generated by C_4 -group and its monodromy representations are related to each other by the following construction. Take the manifold H/Γ and consider its compactification $\overline{H/\Gamma}$. The genus of $\overline{H/\Gamma}$ can be computed using the Hurwitz's formula. Each closed path γ on $\overline{H/\Gamma}$ corresponds to some closed B_3 - or C_4 -braid. The path can be expanded through the (non-trivial) conjugacy classes of Γ which in turn determine the group $C_4(CP^1)$ due to the uniqueness of the monodromy representation.

In each conjugacy class of the group Γ lies the shortest closed path (geodesics). Since the length of geodesics are completely determined by the spectrum of the Laplacian Δ

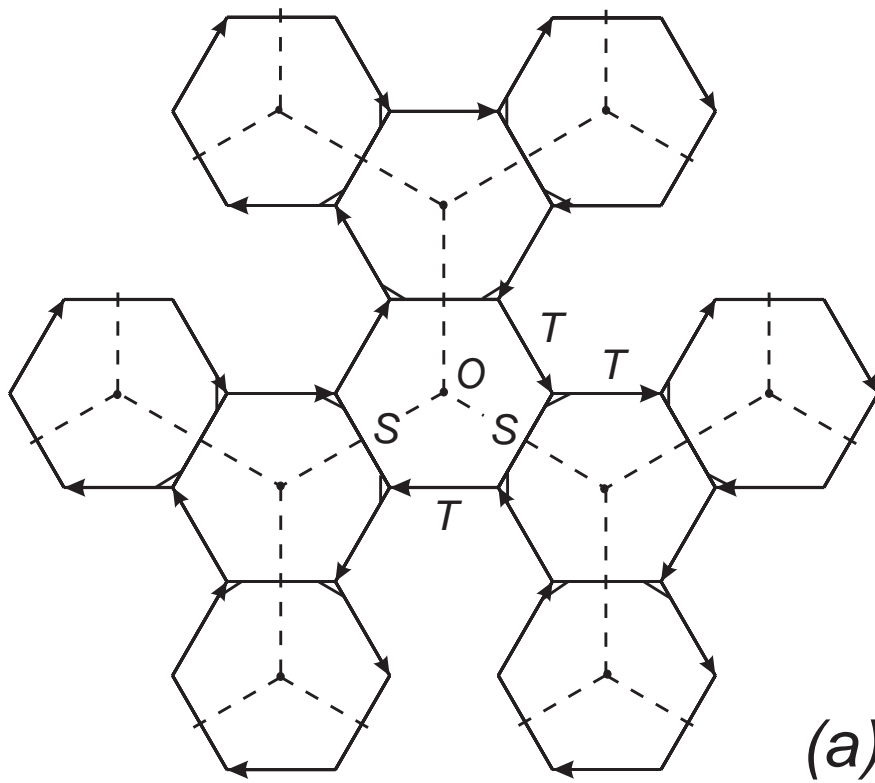
(by the "partition function") on the corresponding Riemann surface, the connection between the knot invariants and the partition function of Δ should exist. The natural formula, linking the spectral properties of the Laplacian on $\overline{H/\Gamma}$ with the conjugacy classes (of hyperbolic, elliptic and parabolic types) of Γ is the Selberg trace formula. Hence, we claim it is possible to express the invariants of knots represented by B_3 - C_4 -braids via Selberg trace formula. More details will be discussed in the forthcoming publication.

Acknowledgements

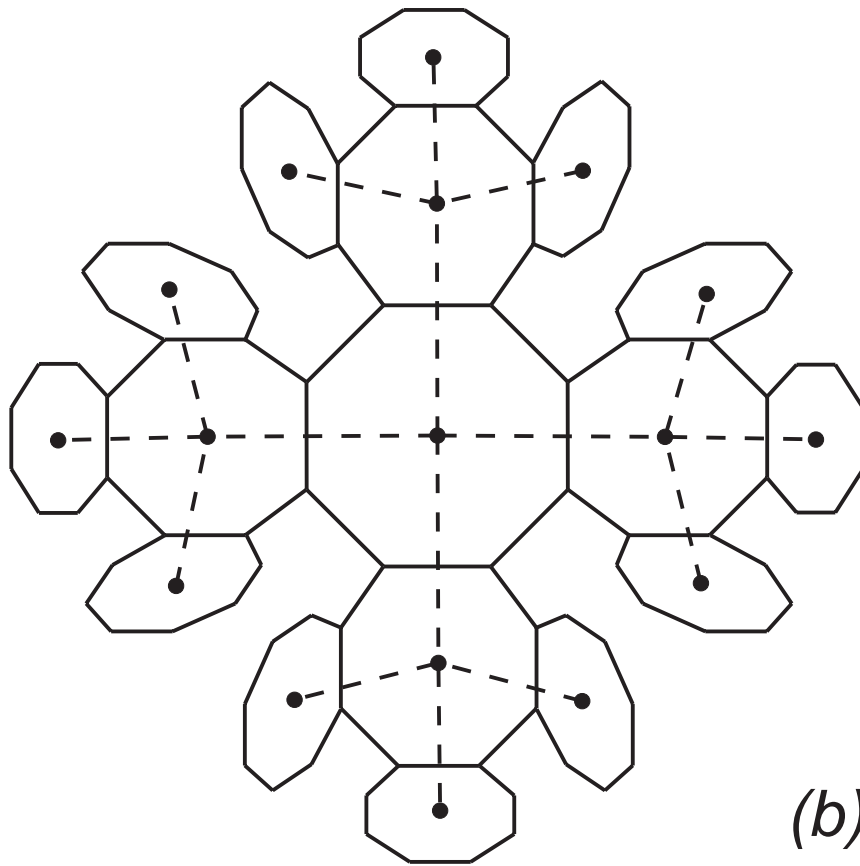
We thank L. Chayes and S. Hirson for the help in final preparation of the paper. This research is partially supported by the RFBR grant 96-01-00744.

References

- [1] F. Wilczek, *Fractional Statistics and Anyon Superconductivity* (WSPC: Singapore, 1990)
- [2] M. Monastyrsky, *Topology of Gauge Fields and Condensed Matter* (Plenum: NY, 1993)
- [3] S. Nechaev, *Statistics of Knots and Entangled Random Walks* (WSPC: Singapore, 1996)
- [4] A.Yu. Grosberg, A.R. Khokhlov, *Statistical Physics of Macromolecules* (AIP Press: NY, 1994)
- [5] *Integrable Models and Strings*, Lect. Notes Phys., 436, (Springer: Heidelberg, 1994)
- [6] E. Witten, *Comm. Math. Phys.*, 121 (1989), 351
- [7] J. Birman, *Knots, Links and Mapping Class Groups*, *Ann. Math. Stud.*, 82, (Princeton Univ. Press: Princeton, NJ, 1976)
- [8] A.B. Venkov, Preprint LOMI P/2/86
- [9] D. V. Anosov, A.A. Bolibruch, *Riemann–Hilbert problem*, (Vieweg: Braunschweig, 1994)
- [10] S. Nechaev, A. Vershik, *J. Phys. (A): Math. & Gen.*, 27 (1994), 2289
- [11] A.A. Belavin, A.M. Polyakov, A.B. Zamolodchikov, *Nucl. Phys.(B)*, 241 (1984), 333
- [12] V.I.S.Dotsenko, *Nucl.Phys. B-235 [FS11]* (1984), 54
- [13] G.D. Mostov, *Bull. AMS*, 16 (1987), 225
- [14] J. Gilman, *Memoirs AMS*, 1996 (in press)
- [15] V.F.R. Jones, *Subfactors and Knots*, CBMS No. 80, (AMS: Providence, 1991)



(a)



(b)